

Introduction to the post-adiabatic expansion for EMRIs

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Fundamental Physics Meets Waveforms With LISA

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What is the post-adiabatic expansion?

These are all basically the same thing:

- post-adiabatic approximation
- two-timescale expansion
- multiscale expansion
- multi-timescale analysis
- method of averaging
- near-identity averaging transformations

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Methods of separating secular from oscillatory effects

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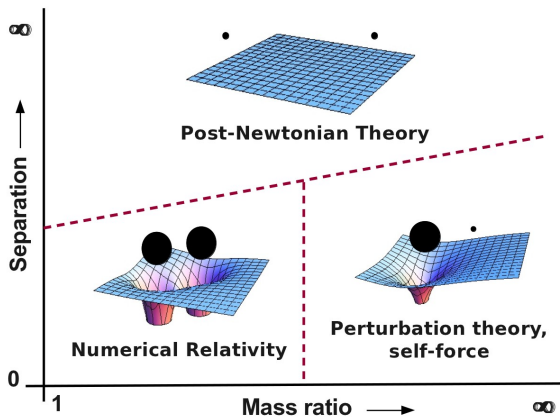
Accurate: uniform accuracy on long and short time scales

Fast: 10 – 100ms per full waveform

Modular: readily accommodate non-GR effects

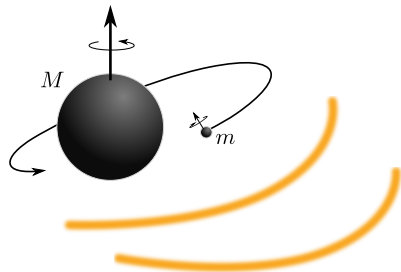
What is the post-adiabatic expansion?

A way of organizing the self-force expansion for extreme mass ratios



[Image credit: Leor Barack]

Gravitational self-force theory



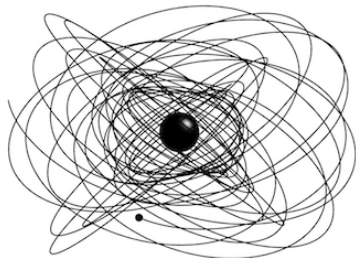
- $\epsilon = 1/q = m/M \ll 1$
- small body perturbs spacetime:

$$g_{\mu\nu} = g_{\mu\nu}^{\text{Kerr}} + \epsilon h_{\mu\nu}^{(1)} + \epsilon^2 h_{\mu\nu}^{(2)} + \dots$$

- perturbation affects m 's motion:

$$\frac{D^2 z^\mu}{d\tau^2} = \epsilon f_{(1)}^\mu + \epsilon^2 f_{(2)}^\mu + \dots$$

Zeroth order: test mass on a geodesic in Kerr



[image credit: Steve Drasco]

- constants $J_A = (M, a, E, L_z, Q)$:
 - ① energy E
 - ② angular momentum L_z
 - ③ Carter constant Q , related to orbital inclination
- phases $\varphi_A = (\varphi_r, \varphi_\theta, \varphi_\phi)$ with frequencies $\Omega_A(J_B)$

- simple ODEs:

$$\frac{d\varphi_A}{dt} = \Omega_A(J_B)$$

$$\frac{dJ_A}{dt} = 0$$

Post-adiabatic expansion

[Flanagan, Hinderer, Lynch, Miller, Moxon, AP, van de Meent, Warburton, ...]

- full, evolving system:

$$\begin{aligned}\frac{d\tilde{\varphi}_A}{dt} &= \Omega_A(\tilde{J}_B) \\ \frac{d\tilde{J}_A}{dt} &= \epsilon \left[F_A^{(0)}(\tilde{J}_B) + \epsilon F_A^{(1)}(\tilde{J}_B) + \mathcal{O}(\epsilon^2) \right]\end{aligned}$$

- waveform

$$h_{\ell m} = \sum_{k^i} \left[\epsilon h_{\ell m k^i}^{(1)}(\tilde{J}_A) + \epsilon^2 h_{\ell m k^i}^{(2)}(\tilde{J}_A) + \mathcal{O}(\epsilon^3) \right] e^{-i(m\tilde{\varphi}^\phi + k^i \tilde{\varphi}_i)}$$

- secondary spin: (i) new slow parameters
(ii) new precession phase

Offline step

- solve field equations for amplitudes $h_{\ell m k^i}^{(n)}$ and forcing functions $F_A^{(n-1)}$ on grid of \tilde{J}_A values

Online step

- solve ODEs for $\tilde{\varphi}_A$ and \tilde{J}_A , add up mode amplitudes $h_{\ell m k^i}^{(n)}$
- FastEMRIWaveforms (**FEW**) software package:
waveform in $\sim 10 - 100\text{ms}$
[Chapman-Bird, Chua, Hughes, Katz, Speri, Warburton, ...]

- phases have an expansion

$$\tilde{\varphi}_A = \epsilon^{-1} \varphi_A^{(0)}(\epsilon t) + \epsilon^0 \varphi_A^{(1)}(\epsilon t) + O(\epsilon)$$

- a model that gets $\varphi_A^{(0)}$ and $\varphi_A^{(1)}$ right should be enough for precise parameter extraction [Burke et al. 2023]
- also need to account for transient orbital resonances (0.5PA)

Adiabatic order (0PA)

determined by $F_A^{(0)}$

- dissipative piece of $f_{(1)}^\mu$

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Accuracy and post-adiabatic counting [Hinderer & Flanagan]

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First post-adiabatic order (1PA)

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- dissipative piece of $f_{(2)}^\mu$
- conservative piece of $f_{(1)}^\mu$
- linear secondary spin effects

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Current state of play

Goal: 1PA waveforms over full parameter space

Background Spacetime	Orbital Configuration	Adiabatic		Post-1-adiabatic		
		1SF (Dissipative)	1SF (Conservative)	2SF (Dissipative)	Spin Effects (Conservative)	Spin Effects (Dissipative)
Schwarzschild	Circular	✓✓✓	✓✓✓	✓✓✓	✓✓✓	✓✓✓
	Eccentric	✓✓✓	✓✓✓	X	✓✓, ✓✓✓*	✓, ✓✓*
Kerr	Circular	✓✓✓	✓✓	X	✓, ✓✓*	✓✓✓*
	Eccentric Equatorial	✓✓✓	✓✓	X	✓, ✓✓*	✓✓*
	Generic	✓✓✓	✓	X	✓	✓*
	Resonances	✓✓✓	✓	X	X	X

✓✓✓ Evolving Waveform ✓✓ Driven Inspiral ✓ Snapshot Calculation *(Anti-)Aligned Spin Only

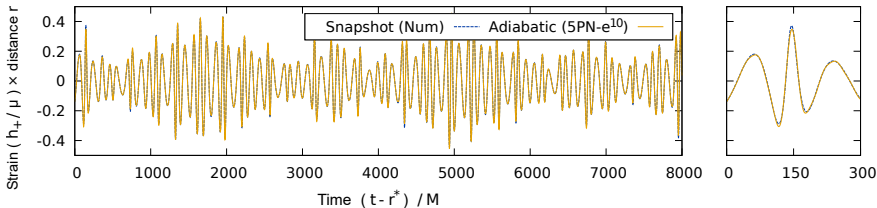
[table courtesy of Josh Mathews]

State of the art in parameter space coverage: 0PA waveforms in Kerr [Hughes et al., Kyoto co., & FEW]

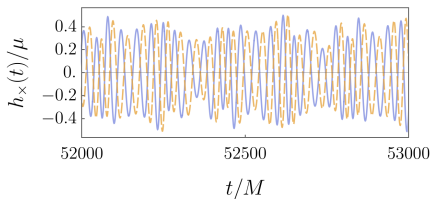
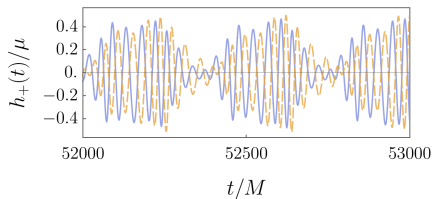
In FEW:

- generic orbits in Kerr:
5.5PN- e^{10} approximation
[Isoyama, Kyoto co., Chua, AP]
- equatorial orbits in Kerr:
fully relativistic waveforms

[animation credit: P. Lynch]



1PA secondary spin effects [Prague groups; Drummond & Hughes et al.]



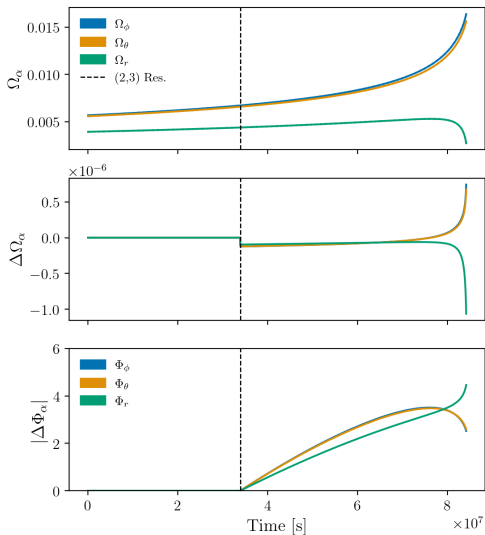
- waveforms including 1PA conservative spin effects “available” for generic configurations in Kerr
- currently missing efficient method of computing non-flux dissipative effects (but see [Grant; Skoupý and Witzany; Mathews et al.])

Resonances [Flanagan & Hinderer, Lynch et al., Isoyama et al., ...]

Ω_r/Ω_θ becomes momentarily rational

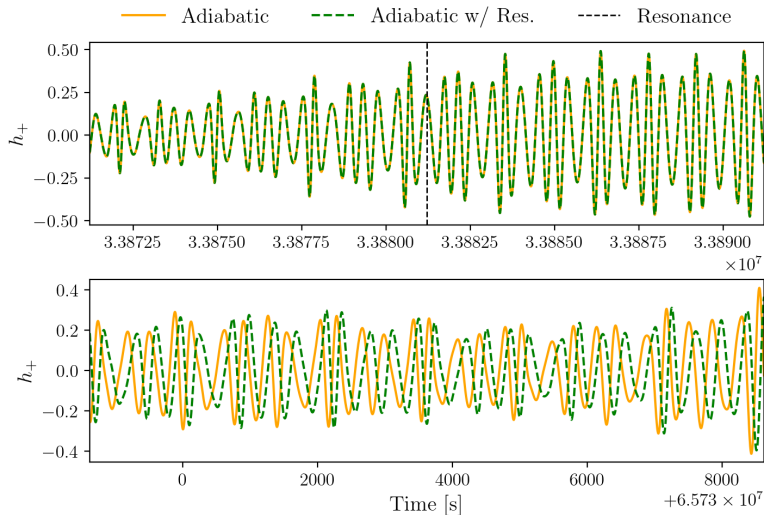
Ω_A jump slightly across the resonance

leads to a significant cumulative change in phase



[Image credit: C. Chapman-Bird]

Resonances in FEW [Chapman-Bird & Warburton]

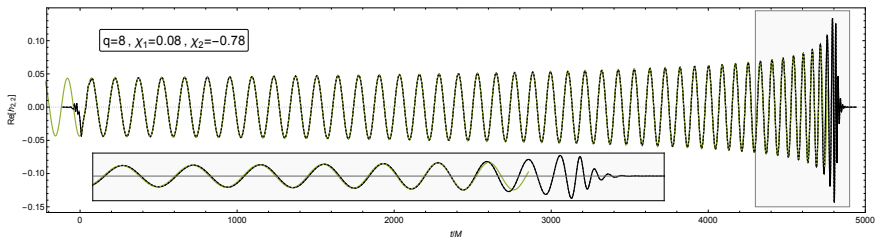


[Image credit: C. Chapman-Bird]

State of the art in accuracy:

Complete 1PA waveforms [Mathews, AP, Warburton, Wardell & MSF group]

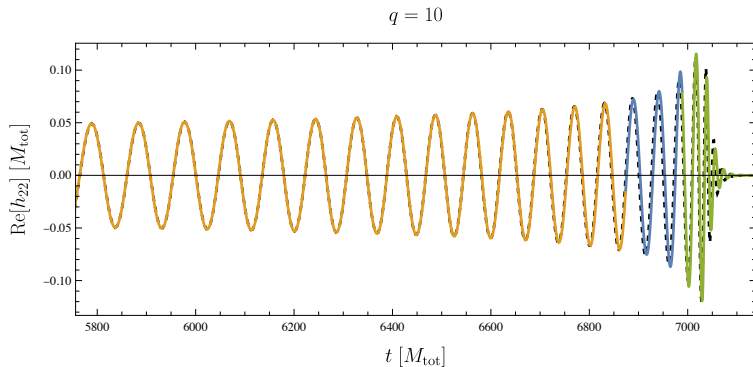
Comparison with NR waveform from SXS collaboration



- complete quasicircular 1PA model with generic secondary spin and linear primary spin
- in FEW: version with zero primary spin, aligned secondary spin

Complete inspiral-merger ringdown models

[Kuchler, Compere, Durkan, AP & MSF group]



- 1PA inspiral, next-to-leading transition, leading-order plunge
- see L. Kuchler's talk

Post-adiabatic self-force waveforms

- high accuracy for IMRIs and EMRIs
- native fast waveform generation
—fast enough for data analysis (?)

Status

- 0PA: equatorial inspirals into Kerr primary
- resonances, secondary spin, precession in the pipeline
- complete 1PA: quasicircular inspiral into slowly spinning primary

Challenge

- coverage of parameter space
- non-GR effects