

On supersymmetry breaking in supergravity

Alessandro Tomasiello

AEI, 3.12.2019

Introduction

Non-supersymmetric solutions of string theory

- have no compelling geometrical story to tell, so far
(unlike susy: generalized/exceptional geometry...)
- for AdS, have been conjectured to be unstable
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Today I'll try to address the first two issues

looking at some large classes of examples.

Plan

- Breaking supersymmetry with pure spinors
- Non-BPS AdS₇ solutions
 - Stability
 - perturbative: part of the KK tower
 - non-perturbative: bubbles

Pure spinors

‘Pure spinor equations’ for BPS solutions

[Graña, Minasian, Petrini, AT ‘05]

$$d_H(e^{3A-\phi}\Phi_+) = 0$$

$$d_H(e^{2A-\phi}\text{Re}\Phi_-) = 0$$

$$d_H(e^{4A-\phi}\text{Im}\Phi_-) = e^{4A} \star \lambda(F)$$

NSNS flux

RR flux

Φ_{\pm} : even/odd ‘pure’ forms

+ algebraic comp. condition

metric determined from Φ_{\pm}

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An easy class:

$$ds^2 = S^{-1/2} ds_{\text{Mink}_6}^2 + K(S^{-1/2} dz^2 + S^{1/2} ds_{\mathbb{R}^3}^2)$$

[motivated by NS5-D6-D8]

[Imamura ‘01;
Janssen, Meessen, Ortin ‘99]

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take Φ_{\pm} adapted to this metric

↓

fluxes are determined. Bianchi ⇒

$$\Delta_3 S + \frac{1}{2} \partial_z^2 S^2 = 0$$

$$K = -\frac{4}{F_0} \partial_z S$$

rare case where **single PDE**

$S(x_1, x_2, x_3, z)$
[Legramandi, AT ‘19]

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- $S(x_1, z)$ ‘hodograph’ transf. **linearizes it:**

$$\begin{array}{l} x_1 \equiv \partial_U V \\ z \equiv \partial_S V \end{array} \Rightarrow \partial_S^2 V + S \partial_U^2 V = 0 \quad \text{‘Tricomi equation’}$$

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↓
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- $S = kz + S_3(x) \quad \Delta_3 S_3 = -k^2$ gen. Green function; **analytic** solution

can be made **compact** with $O8_{\pm}$

susy breaking:

keep same metric & RR fluxes; change NSNS
impose Bianchi, but not BPS

[Legramandi, AT '19]

$$\Delta_3 S + \frac{1}{2} \partial_z^2 S^2 = 0 \rightsquigarrow \Delta_3 S + \frac{1}{2} \partial_z^2 S^2 + c(c - 2\partial_z S) = 0$$

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- we have worked out several other classes

taken or generalized from
classif. of $\text{Mink}_4 \times S^2$ solutions

[Macpherson, T '17; Apruzzi, Geipel, Legramandi,
Macpherson, Zagermann '18; Legramandi, Macpherson '18]

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- Why does this work?

Susy-breaking in AdS₇

- All AdS₇ solutions in type II:

[Apruzzi, Fazzi, Rosa, AT '13
Apruzzi, Fazzi, Passias, Rota, AT '15;
Cremonesi, AT '15]

$$\frac{1}{\pi\sqrt{2}}ds^2 = 8\sqrt{-\frac{\alpha}{\ddot{\alpha}}}ds_{\text{AdS}_7}^2 + \sqrt{-\frac{\ddot{\alpha}}{\alpha}} \left(dz^2 + \frac{\alpha^2}{\dot{\alpha}^2 - 2\alpha\ddot{\alpha}} ds_{S^2}^2 \right)$$

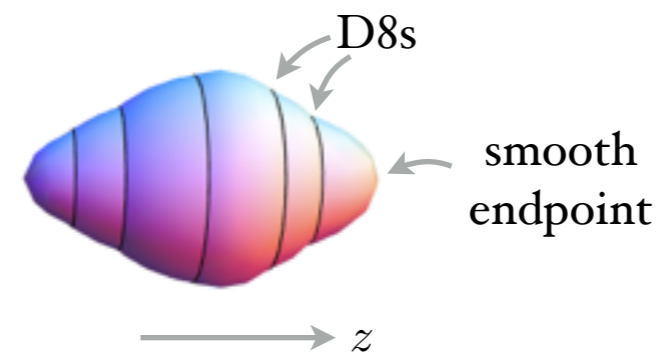
interval

$$\ddot{\alpha} = F_0 \quad \Rightarrow \quad \alpha \text{ piecewise cubic}$$

$$e^\phi = 2^{1/4} 162\pi^{5/2} \frac{(-\alpha/\ddot{\alpha})^{3/4}}{\sqrt{\dot{\alpha}^2 - 2\alpha\ddot{\alpha}}}$$

$$B = \pi \left(-z + \frac{\alpha\dot{\alpha}}{\dot{\alpha}^2 - 2\alpha\ddot{\alpha}} \right) \text{vol}_{S^2}$$

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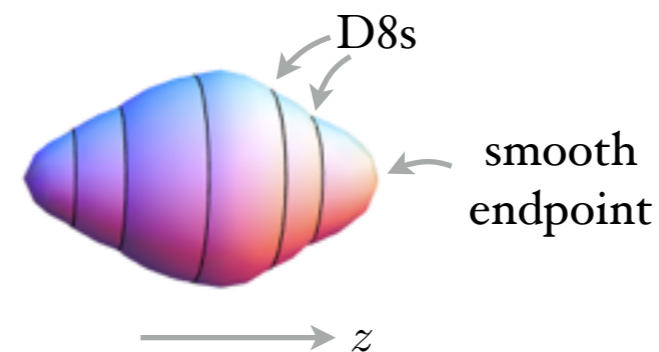
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- position of a D8:

$$z = k = \text{its D6-charge}$$

- $z \in [0, N]$

NS₅ flux



- Any AdS₇ solution in IIA \mapsto **consistent truncation** to ‘minimal gauged 7d sugra’

[Passias, Rota, AT '15;
Malek, Samtleben, Vall Camell '19]

fields: $g_{\mu\nu}^{(7)}$, A_{μ}^i , X

[susy vacuum: $X = 1$]

uplift:

$$\frac{1}{\pi\sqrt{2}} ds^2 = \frac{120}{V(X)\sqrt{X}} \sqrt{-\frac{\alpha}{\ddot{\alpha}}} ds_{\text{AdS}_7}^2 + X^{5/2} \sqrt{-\frac{\ddot{\alpha}}{\alpha}} \left(dz^2 + \frac{\alpha^2}{\dot{\alpha}^2 - 2X^5\alpha\ddot{\alpha}} ds_{S^2}^2 \right)$$

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- AdS₇ \rightarrow AdS _{d} \times Σ_{7-d} comp. & RG flows

- non-susy AdS₇ solution: $X^5 = \frac{1}{2}$

So we get a susy-breaking sister solution:

$$\frac{1}{\pi\sqrt{2}} ds^2 = \sqrt{-\frac{\alpha}{\ddot{\alpha}}} ds_{\text{AdS}_7}^2 + \sqrt{-\frac{\ddot{\alpha}}{\alpha}} \left(dz^2 + \frac{\alpha^2}{\dot{\alpha}^2 - \alpha\ddot{\alpha}} ds_{S^2}^2 \right)$$

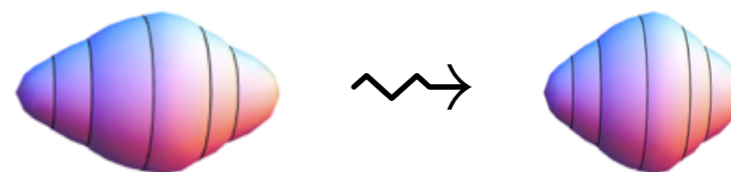
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- Can we have a 7d theory with more than one susy vacuum?

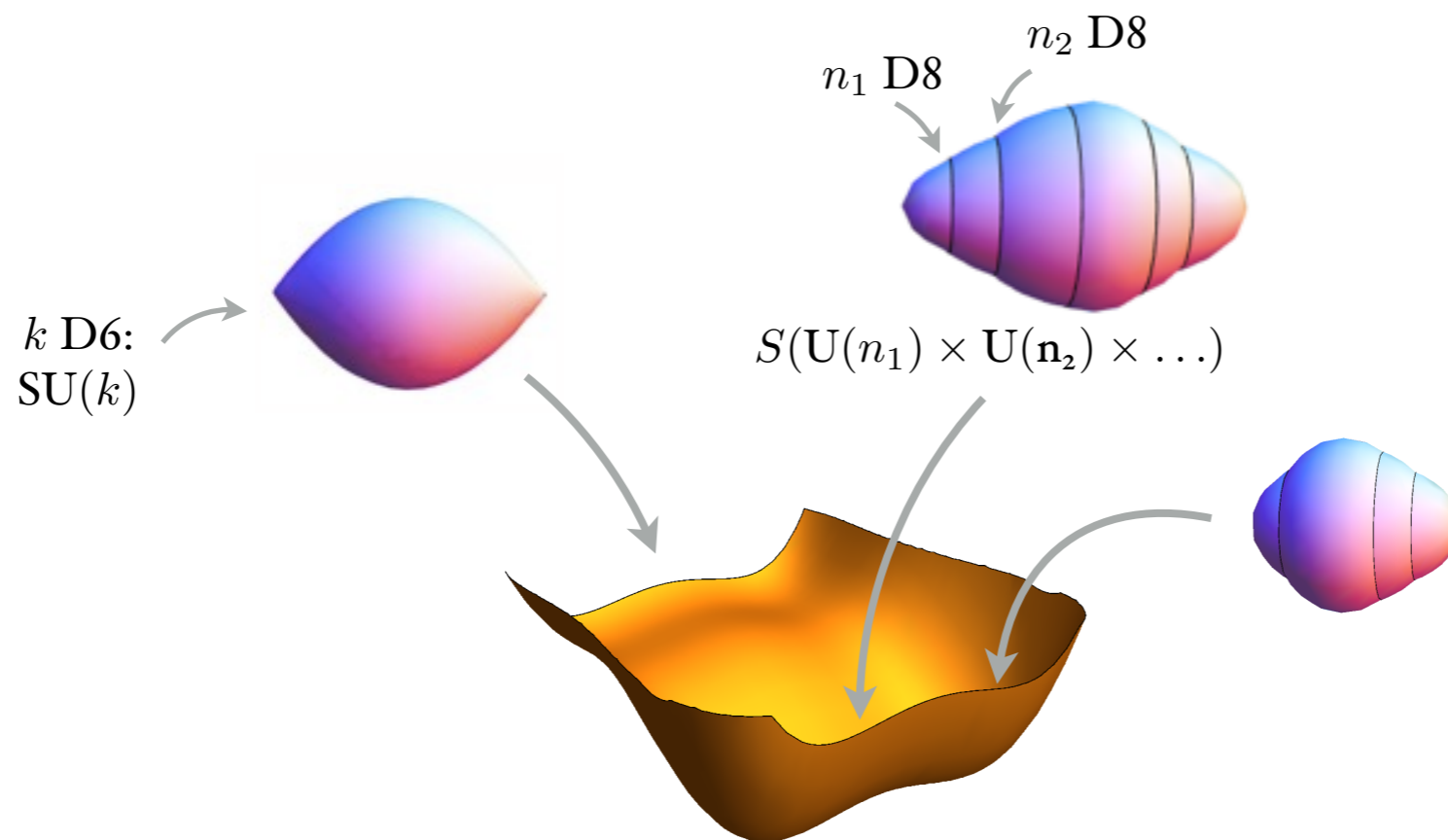
couple minimal 7d gauged sugra to $SU(k)^2$ vector multiplets

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couple minimal 7d gauged sugra to $SU(k)^2$ vector multiplets [De Luca, Gnechhi, Lo Monaco, AT '18]

- many nonabelian vacua \cong AdS₇ vacua with different D8 configurations
- susy RG flows \rightarrow Nahm equations. Agree with CFT expectations



AdS₇ (in)stability

Possible instabilities:

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- ‘Bubbles’. Via tunnel effect, bubble of true vacuum develops and then grows classically

Supersymmetric vacua are protected against both.

But non-supersymmetric ones?

[Breitenlohner, Freedman '82,
Gibbons, Hull, Warner '83...]

- ‘Tachyonic’ modes: Full KK computation is challenging.

- Spin-two masses: simple universal equation

[Bachas, Estes '11]

spectrum of ‘warped Laplacian’

$$\frac{e^{-5A+2\phi}}{\sqrt{\hat{g}}} \partial_m \left(e^{7A-2\phi} \sqrt{\hat{g}} \hat{g}^{mn} \partial_n \right)$$

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Universal bound:

realized for $\psi = \alpha^{-l} Y_l$

[Passias, AT '16]

↑ spherical harm. on S^2

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- **supersymmetric** case: $m^2 \geq 8l(2l + 3) \geq 40$

[in units where $L_{\text{AdS}} = 1$]

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⇒ no separation of scales

proof for susy AdS₇!

[notice: reasonable est. $M_{\text{KK}} \sim \int e^{4A} R_6$ can fail for solutions with O-planes]

- Spin-zero: we can compute many masses in 7d gauged sugra

- also: nonabelian DBI \Rightarrow same results

[Apruzzi, De Luca, Gnecci, Lo Monaco, AT to appear]

potential: $-4\Phi^i\Phi^i - 2\epsilon_{ijk}[\Phi_i, \Phi_j]\Phi_k + [\Phi_i, \Phi_j][\Phi^i, \Phi^j]$

vacua: $[\Phi^i, \Phi^j] = \epsilon^{ijk}\Phi^k$

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- susy-breaking case:

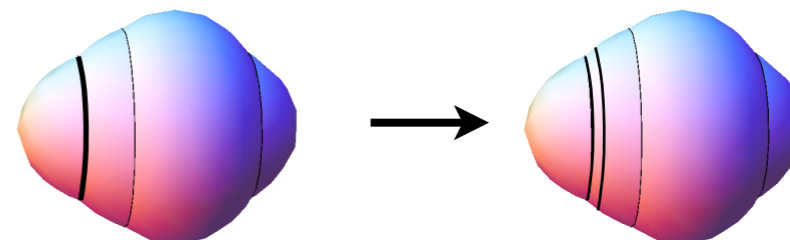
[Apruzzi, De Luca, Gnecci, Lo Monaco, AT to appear]

$$m^2 = 3(k+2)(k-2) = -12, -9, 0, 15, 36, \dots$$

\uparrow
below the BF bound!

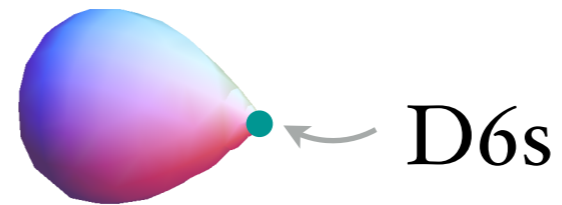
but this -12 appears only with coincident D8s

plausible
scenario:



- Another pert. instability exists for solutions with D6s

such as the non-susy
version of the ‘teardrop’



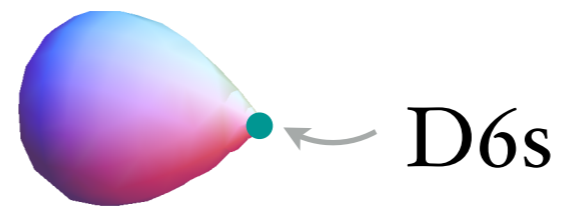
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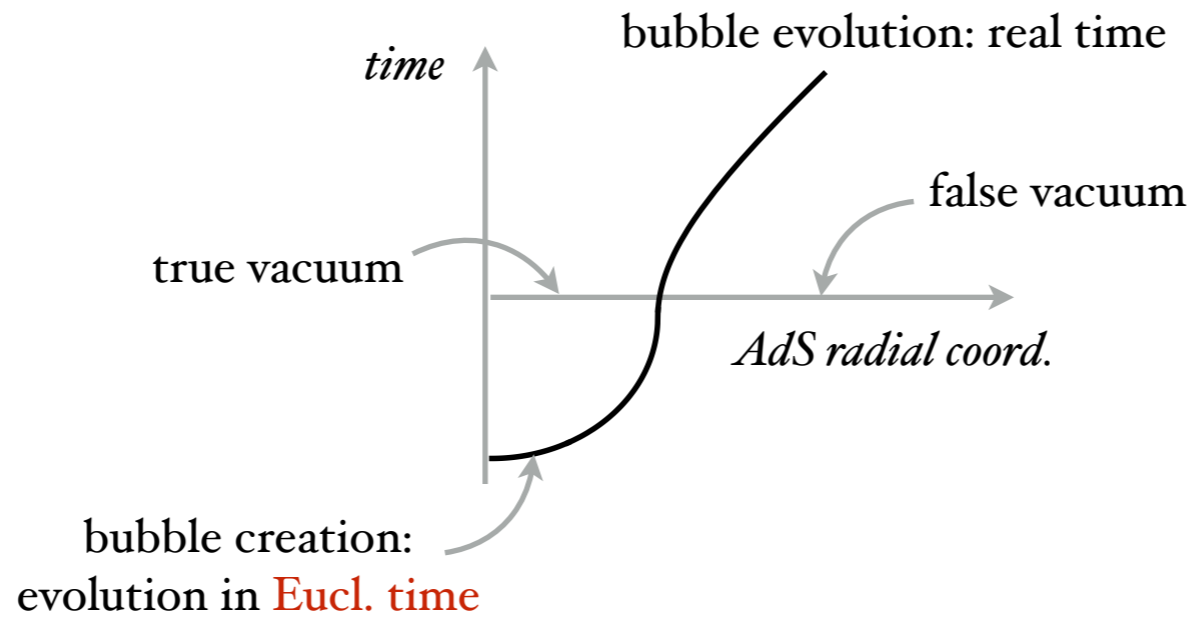
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endpoint: D6 smeared solution?

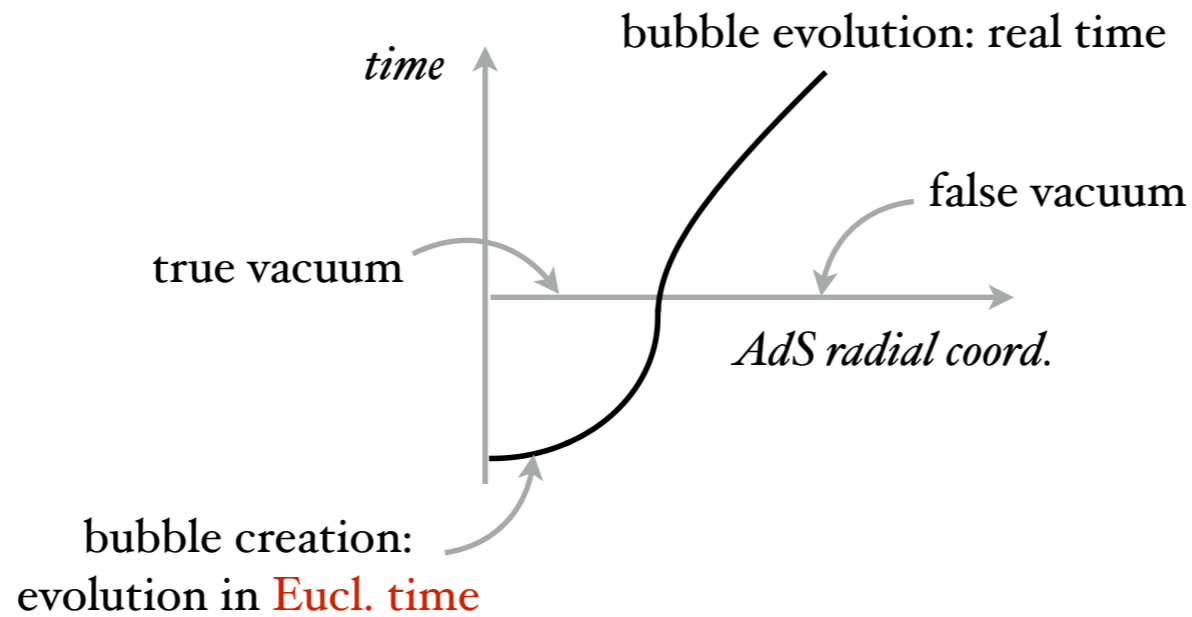
[Blåbäck, Danielsson, Junghans,
Van Riet, Wrase, Zagermann '11]

- ‘Bubbles’.



[Coleman, De Luccia '80]
[Maldacena, Michelson, Strominger '98]

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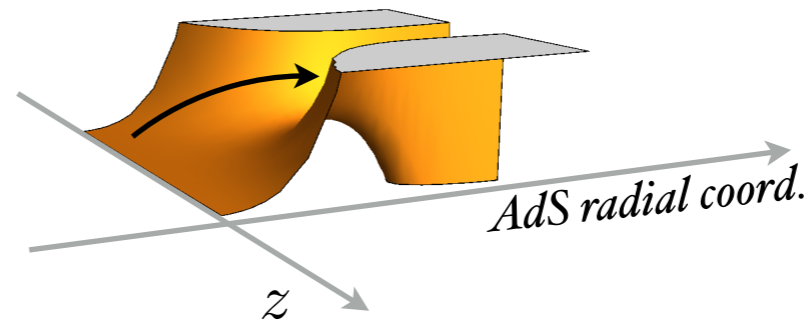


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- NS5 bubble.

tunnel effect: NS5 gets created

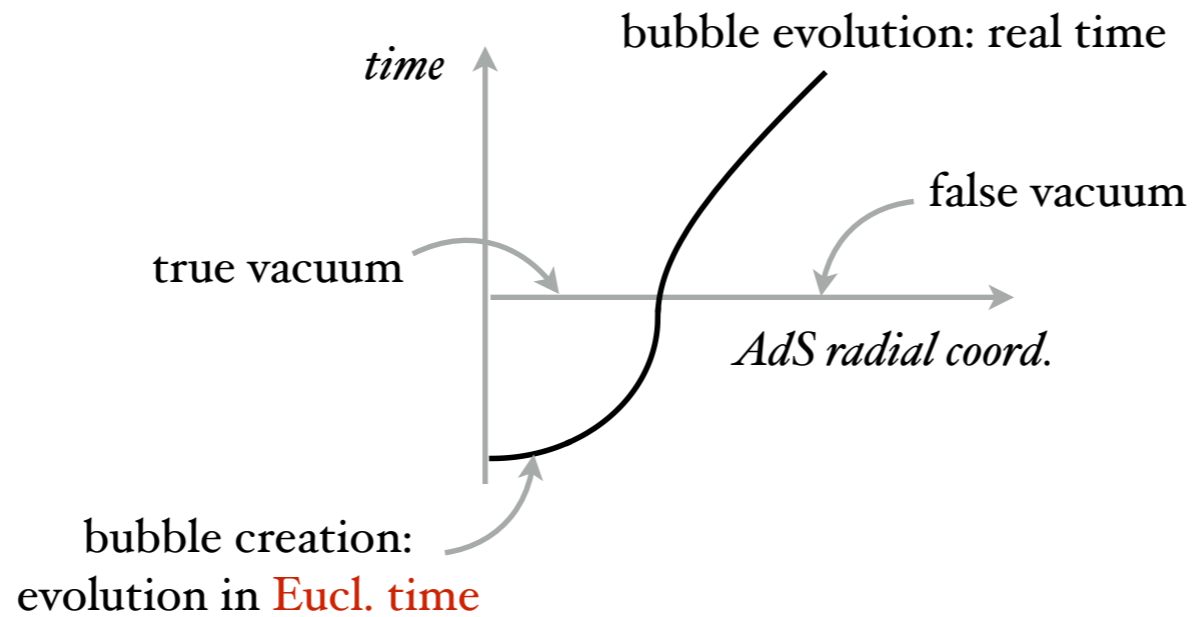
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- this is the type of process that led to the instability conjecture

[Ooguri, Vafa '16,
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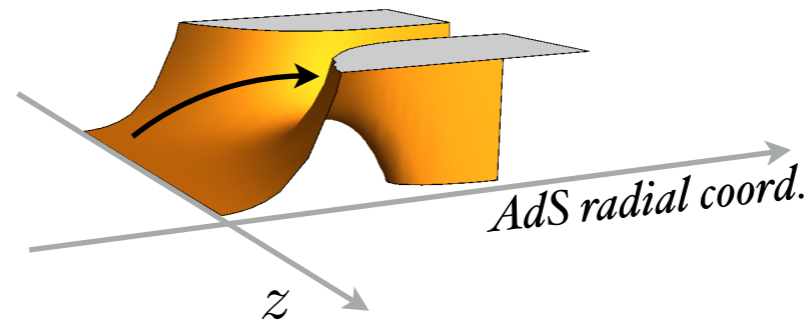


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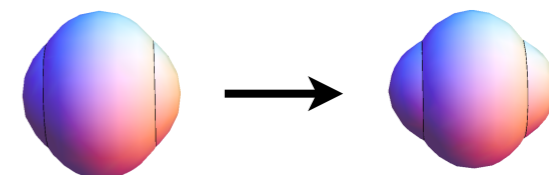
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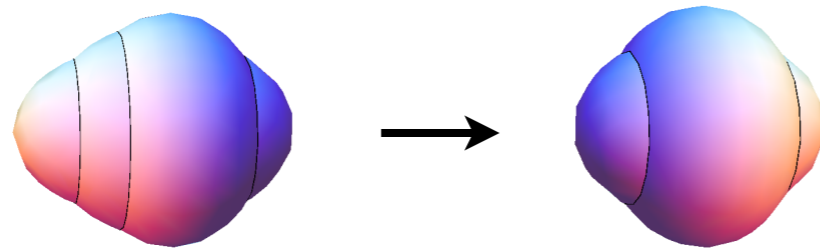
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- this process ‘shortens’ the region where $F_0 = 0$



- Other possible tunnel effect:

[motivated by nonabelian potential]



2 D8s combine: one of them keeps all D6 charge, the other tunnels away

- Together these effects seem to rule everything out.

[Apruzzi, De Luca, Gnecci, Lo Monaco, AT to appear]

Conclusions

- Steps towards procedure to break supersymmetry
in the Minkowski case [Legramandi, AT '19]
- Older method: consistent truncation. Are non-susy AdS₇ stable?
- Tachyons make D8s repel each other
- Tunnel effects probably make remaining ones unstable