

Conformal Standard Model

Krzysztof A. Meissner
University of Warsaw
AEI Potsdam

HermannFest, AEI, 6.09.2012

Hierarchy Problem

- Standard Model is conformally invariant **except** for the tree level Higgs mass term

$$-m^2\Phi^\dagger\Phi$$

Hierarchy Problem

- Standard Model is conformally invariant **except** for the tree level Higgs mass term

$$-m^2\Phi^\dagger\Phi$$

- Quantum corrections $m^2 \sim \Lambda^2 \Rightarrow$
why $m \ll M_P$?
(with UV cutoff $\Lambda =$ scale of ‘new physics’)

Hierarchy Problem

- Standard Model is conformally invariant **except** for the tree level Higgs mass term
$$-m^2\Phi^\dagger\Phi$$
- Quantum corrections $m^2 \sim \Lambda^2 \Rightarrow$
why $m \ll M_P$?
(with UV cutoff $\Lambda =$ scale of ‘new physics’)
- Most popular proposal: SM \longrightarrow (N)MSSM
(but LHC does not seem to pointedly like it...)

Hierarchy Problem

- Standard Model is conformally invariant **except** for the tree level Higgs mass term
$$-m^2\Phi^\dagger\Phi$$
- Quantum corrections $m^2 \sim \Lambda^2 \Rightarrow$
why $m \ll M_P$?
(with UV cutoff $\Lambda =$ scale of ‘new physics’)
- Most popular proposal: SM \longrightarrow (N)MSSM
(but LHC does not seem to pointedly like it...)
- conformal symmetry also solves the HP:
only 2 new parameters (MSSM 116)
built-in breaking (conf. anomaly,
Coleman-Weinberg)

Conformal Standard Model – assumptions

K.A.M., H. Nicolai, Phys.Lett. B648 (2007) 312

- **Classical conformal symmetry** in SM, all masses from conformal anomaly

Conformal Standard Model – assumptions

K.A.M., H. Nicolai, Phys.Lett. B648 (2007) 312

- **Classical conformal symmetry** in SM, all masses from conformal anomaly
- see-saw mechanism for neutrinos

Conformal Standard Model – assumptions

K.A.M., H. Nicolai, Phys.Lett. B648 (2007) 312

- **Classical conformal symmetry** in SM, all masses from conformal anomaly
- see-saw mechanism for neutrinos
- the model viable up to M_{Pl} , no GUTs, no new scales, no low energy SUSY

Conformal Standard Model – assumptions

K.A.M., H. Nicolai, Phys.Lett. B648 (2007) 312

- **Classical conformal symmetry** in SM, all masses from conformal anomaly
- see-saw mechanism for neutrinos
- the model viable up to M_{Pl} , no GUTs, no new scales, no low energy SUSY
- classical conformal symmetry \Rightarrow Ward Ids \Rightarrow all corrections logarithmic (Bardeen)

K.A.M., H. Nicolai, Phys. Lett. B660 (2008) 260

Nonconformal \rightarrow conformal?

- we expect at M_{Pl} a non-conformal theory
is it possible to get at low energies
a (classically) conformal FT + gravity?

Nonconformal \rightarrow conformal?

- we expect at M_{Pl} a non-conformal theory
is it possible to get at low energies
a (classically) conformal FT + gravity?
- example:
gauged $N = 4$ supergravity with $\kappa \rightarrow 0$
 \Rightarrow supergravity + conformal $N = 4$ SYM

K.A.M., H. Nicolai, Phys.Rev. D80:086005,2009

The Model

- Conformally invariant $\mathcal{L} = \mathcal{L}_{kin} + \mathcal{L}'$:

$$\begin{aligned} \mathcal{L}' := & \left(\bar{L}^i \Phi Y_{ij}^E E^j + \bar{Q}^i \epsilon \Phi^* Y_{ij}^D D^j + \bar{Q}^i \epsilon \Phi^* Y_{ij}^U U^j + \right. \\ & \left. + \bar{L}^i \epsilon \Phi^* Y_{ij}^\nu \nu_R^j + \varphi \nu_R^{iT} C Y_{ij}^M \nu_R^j + \text{h.c.} \right) \\ & - \frac{\lambda_1}{4} (\Phi^\dagger \Phi)^2 - \frac{\lambda_2}{2} |\varphi|^2 (\Phi^\dagger \Phi) - \frac{\lambda_3}{4} |\varphi|^4 \end{aligned}$$

The Model

- Conformally invariant $\mathcal{L} = \mathcal{L}_{kin} + \mathcal{L}'$:

$$\begin{aligned} \mathcal{L}' := & \left(\bar{L}^i \Phi Y_{ij}^E E^j + \bar{Q}^i \epsilon \Phi^* Y_{ij}^D D^j + \bar{Q}^i \epsilon \Phi^* Y_{ij}^U U^j + \right. \\ & \left. + \bar{L}^i \epsilon \Phi^* Y_{ij}^\nu \nu_R^j + \varphi \nu_R^{iT} \mathcal{C} Y_{ij}^M \nu_R^j + \text{h.c.} \right) \\ & - \frac{\lambda_1}{4} (\Phi^\dagger \Phi)^2 - \frac{\lambda_2}{2} |\varphi|^2 (\Phi^\dagger \Phi) - \frac{\lambda_3}{4} |\varphi|^4 \end{aligned}$$

- $Y_{ij}^M \sim O(1)$, $Y_{ij}^\nu \sim O(10^{-6})$ (see-saw)

The Model

- Conformally invariant $\mathcal{L} = \mathcal{L}_{kin} + \mathcal{L}'$:

$$\begin{aligned} \mathcal{L}' := & \left(\bar{L}^i \Phi Y_{ij}^E E^j + \bar{Q}^i \epsilon \Phi^* Y_{ij}^D D^j + \bar{Q}^i \epsilon \Phi^* Y_{ij}^U U^j + \right. \\ & \left. + \bar{L}^i \epsilon \Phi^* Y_{ij}^\nu \nu_R^j + \varphi \nu_R^{iT} \mathcal{C} Y_{ij}^M \nu_R^j + \text{h.c.} \right) \\ & - \frac{\lambda_1}{4} (\Phi^\dagger \Phi)^2 - \frac{\lambda_2}{2} |\varphi|^2 (\Phi^\dagger \Phi) - \frac{\lambda_3}{4} |\varphi|^4 \end{aligned}$$

- $Y_{ij}^M \sim O(1)$, $Y_{ij}^\nu \sim O(10^{-6})$ (see-saw)
- new (complex) scalar necessary!

The Model

- Standard Model + very light neutrinos
+ heavy neutrinos (~ 1 TeV)
+ complex new scalar

The Model

- Standard Model + very light neutrinos
+ heavy neutrinos (~ 1 TeV)
+ complex new scalar
- Calculation of the CW effective action
– numerical analysis (no 2 loops...)

The Model

- Standard Model + very light neutrinos
+ heavy neutrinos (~ 1 TeV)
+ complex new scalar
- Calculation of the CW effective action
– numerical analysis (no 2 loops...)
- BEH mechanism for EW symm. $\langle \Phi \rangle \neq 0$
SSB of the lepton number symm. $\langle \varphi \rangle \neq 0$

The Model

- Standard Model + very light neutrinos + heavy neutrinos (~ 1 TeV) + complex new scalar
- Calculation of the CW effective action – numerical analysis (no 2 loops...)
- BEH mechanism for EW symm. $\langle \Phi \rangle \neq 0$
SSB of the lepton number symm. $\langle \varphi \rangle \neq 0$
- axion – (pseudo)GB of global lepton number symmetry

$$L^i \rightarrow e^{i\alpha} L^i, \quad E^i \rightarrow e^{i\alpha} E^i, \quad \nu_R^i \rightarrow e^{i\alpha} \nu_R^i, \quad \varphi \rightarrow e^{-2i\alpha} \varphi$$

Results

- At the minimum Higgs and the new scalar mix

$$H' = H \cos \beta + \phi \sin \theta, \quad \phi' = -H \sin \beta + \phi \cos \theta$$

of masses M_1 and M_2 .

K.A.M., H. Nicolai, Phys.Lett. B648 (2007) 312

Results

- At the minimum Higgs and the new scalar mix

$$H' = H \cos \beta + \phi \sin \theta, \quad \phi' = -H \sin \beta + \phi \cos \theta$$

of masses M_1 and M_2 .

K.A.M., H. Nicolai, Phys.Lett. B648 (2007) 312

- propagator

$$\frac{i \cos^2 \theta}{p^2 - M_1^2 + iM_1 \Gamma_{SM}(M_1) \cos^2 \theta} + \left\{ \begin{array}{l} \cos \theta \rightarrow \sin \theta \\ M_1 \rightarrow M_2 \end{array} \right\}$$

in cross section there should be two “bumps”:
heights as for Higgs
widths smaller by $\cos^2(\sin^2)\theta$

Phenomenology

- no other new particles at LHC except standard Higgs and new scalar

Phenomenology

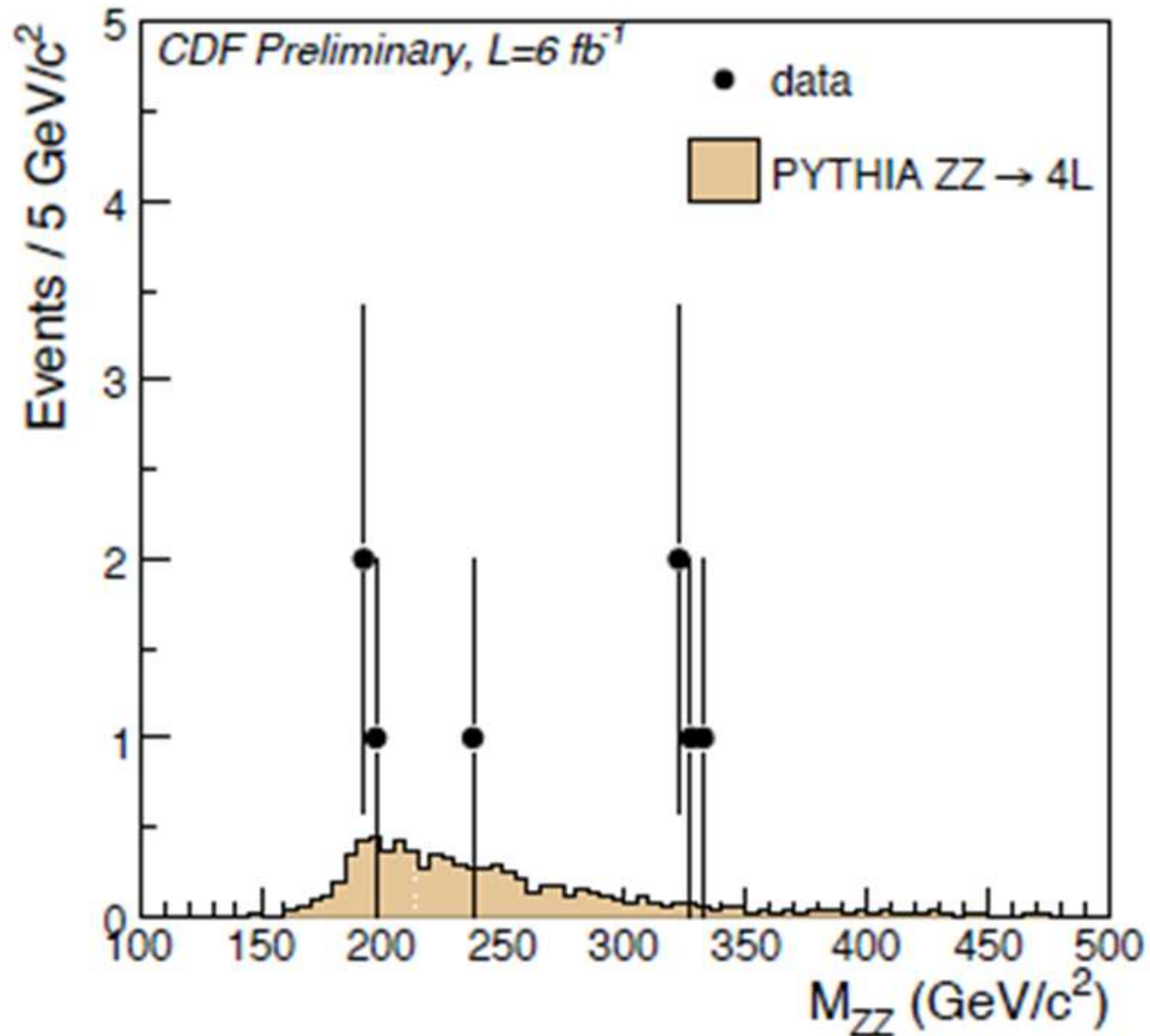
- no other new particles at LHC except standard Higgs and new scalar
- LHC has announced the mass 125 GeV ($M_1?$)

Phenomenology

- no other new particles at LHC except standard Higgs and new scalar
- LHC has announced the mass 125 GeV (M_1 ?)
- CDF has seen a 5σ excess of events $2Z \rightarrow l^+l^-l^+l^-$ at the same invariant mass 325 GeV we conjecture that it is a narrow resonance corresponding to ϕ' ($M_2 = 325$ GeV?)

K.A.M., H. Nicolai, arXiv:1208.5653[hep-ph]

CDF 4 lepton cases



Axion

K.A.M., H. Nicolai, Eur.Phys.J.C57:493,2008

- axion: phase of the new scalar ϕ
Goldstone boson of spontaneously broken
lepton number symmetry

Axion

K.A.M., H. Nicolai, Eur.Phys.J.C57:493,2008

- axion: phase of the new scalar ϕ
Goldstone boson of spontaneously broken lepton number symmetry
- the $a\gamma\gamma$ coupling is calculable and very small (because of small neutrino masses!)

$$\approx \frac{\alpha_w \alpha m_\nu}{8\pi^2 M_W^2} a \vec{E} \cdot \vec{B} \approx 10^{-15} \text{ GeV}^{-1} a \vec{E} \cdot \vec{B}$$

Axion

K.A.M., H. Nicolai, Eur.Phys.J.C57:493,2008

- axion: phase of the new scalar ϕ
Goldstone boson of spontaneously broken lepton number symmetry
- the $a\gamma\gamma$ coupling is calculable and very small (because of small neutrino masses!)

$$\approx \frac{\alpha_w \alpha m_\nu}{8\pi^2 M_W^2} a \vec{E} \cdot \vec{B} \approx 10^{-15} \text{ GeV}^{-1} a \vec{E} \cdot \vec{B}$$

- agg coupling (calculable)
non-local, for small momenta \sim anomaly
 \Rightarrow solution of the strong CP problem

Axion

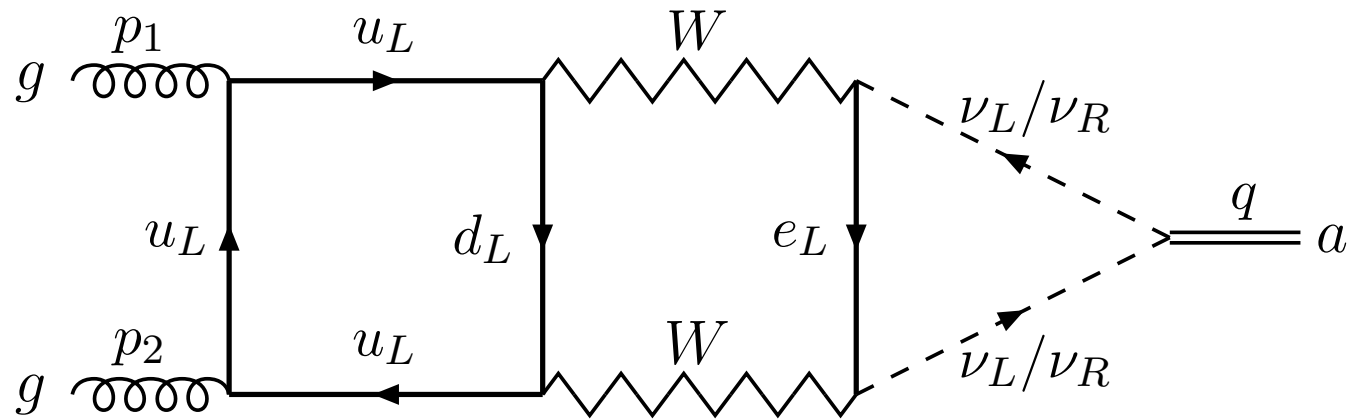
K.A.M., H. Nicolai, *Eur.Phys.J.C*57:493,2008

- axion: phase of the new scalar ϕ
Goldstone boson of spontaneously broken lepton number symmetry
- the $a\gamma\gamma$ coupling is calculable and very small (because of small neutrino masses!)

$$\approx \frac{\alpha_w \alpha m_\nu}{8\pi^2 M_W^2} a \vec{E} \cdot \vec{B} \approx 10^{-15} \text{ GeV}^{-1} a \vec{E} \cdot \vec{B}$$

- agg coupling (calculable)
non-local, for small momenta \sim anomaly
 \Rightarrow solution of the strong CP problem
- a very good candidate for CDM

Axion coupling to gluons



A. Latosiński, K.A.M., H. Nicolai, arXiv:1203.3886 [hep-ph]

Summary

- assumption of classical CS solves problems of SM (hierarchy, strong CP ,...)

Summary

- assumption of classical CS solves problems of SM (hierarchy, strong CP ,...)
- the model has 1 more parameter than SM (MSSM 116) and uses quantum symmetry breaking (conformal anomaly)

Summary

- assumption of classical CS solves problems of SM (hierarchy, strong CP ,...)
- the model has 1 more parameter than SM (MSSM 116) and uses quantum symmetry breaking (conformal anomaly)
- it has a definite (unique) prediction for LHC – 2 scalars with SM Higgs cross sections scaled by $\cos^2(\sin^2)\theta$, no other particles

Summary

- assumption of classical CS solves problems of SM (hierarchy, strong CP ,...)
- the model has 1 more parameter than SM (MSSM 116) and uses quantum symmetry breaking (conformal anomaly)
- it has a definite (unique) prediction for LHC – 2 scalars with SM Higgs cross sections scaled by $\cos^2(\sin^2)\theta$, no other particles
- there is an extremely light and weakly coupled particle (axion)

Summary

- assumption of classical CS solves problems of SM (hierarchy, strong CP ,...)
- the model has 1 more parameter than SM (MSSM 116) and uses quantum symmetry breaking (conformal anomaly)
- it has a definite (unique) prediction for LHC – 2 scalars with SM Higgs cross sections scaled by $\cos^2(\sin^2)\theta$, no other particles
- there is an extremely light and weakly coupled particle (axion)
- if true the CSM is a theory descending directly from the Planck scale

HAPPY BIRTHDAY!