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# $SO(9)$ supergravity\*

conference on the occasion of  
Hermann Nicolai's 60th anniversary  
Golm 2012



Henning Samtleben



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\* in two dimensions

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# $SO(9)$ supergravity\*

[Thomas Ortiz, HS] 2012

[HS, Martin Weidner] 2007

[Hermann Nicolai, HS] 2000

[Hermann Nicolai, Bernard de Wit] 1982

# Domain wall / QFT correspondence

## holography for Dp-branes

[H.J. Boonstra, K. Skenderis, P. Townsend, 1999]

AdS / CFT case

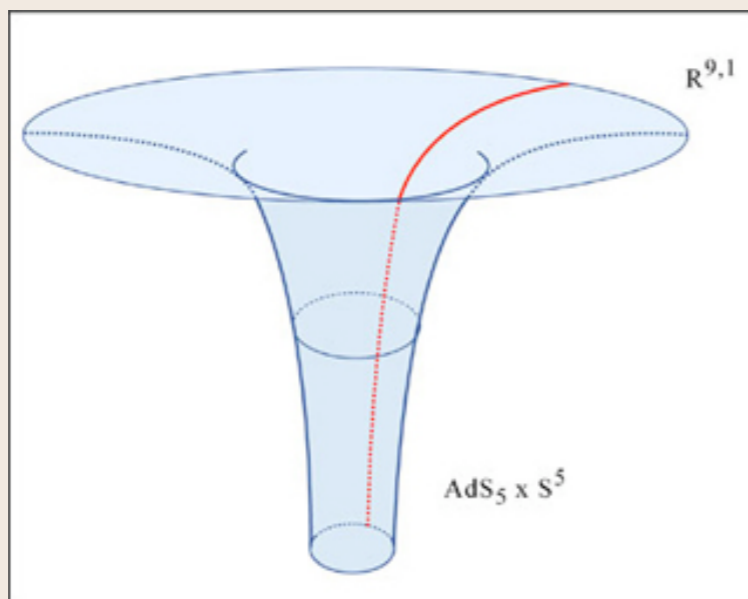
gaugings of maximal supergravity

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D3	IIB	$\text{AdS}_5 \times S^5$	$d=5, \text{SO}(6)$
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[Günaydin, Romans, Warner, 1985]



# Domain wall / QFT correspondence

## holography for Dp-branes

[H.J. Boonstra, K. Skenderis, P. Townsend, 1999]

warped

gaugings of maximal supergravity

D6	IIA	$\text{AdS}_8 \times S^2$	d=8, SO(3)
D5	IIB	$\text{AdS}_7 \times S^3$	d=7, SO(4)
D4	IIA	$\text{AdS}_6 \times S^4$	d=6, SO(5)
D3	IIB	$\text{AdS}_5 \times S^5$	d=5, SO(6)
D2	IIA	$\text{AdS}_4 \times S^6$	d=4, SO(7)
F1/D1	IIA/B	$\text{AdS}_3 \times S^7$	d=3, SO(8)
D0	IIA	$\text{AdS}_2 \times S^8$	d=2, SO(9)

[Salam, Sezgin, 1984]

[Pernici, Pilch, van Nieuwenhuizen, 1984]

[Günaydin, Romans, Warner, 1985]

[Hull, 1984]

[de Wit, Nicolai, 1982]

??? try to construct the SO(9) theory ...

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# 1999 : constructing $SO(9)$ supergravity

[Hermann Nicolai, HS]

# 1999 : trying to construct d=2, SO(9) supergravity

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two-dimensional supergravity is particularly simple :

$$\mathcal{L} = -\frac{1}{4}\sqrt{-g}\rho\left(-R + \text{tr}[P^\mu P_\mu]\right) + \mathcal{L}_{\text{ferm}}(\psi^I, \psi_2^I, \chi^{\dot{A}})$$

coset space sigma model  $E_8/SO(16)$  coupled to dilaton gravity

but has a remarkable structure :

(infinite tower of) dual scalar potentials

$$\partial_\mu Y_M \equiv \varepsilon_{\mu\nu} J_M^\nu$$

conserved  $E_8$  Noether current

etc.

→ classical integrability, affine Lie-Poisson symmetry  $E_9$

off-shell symmetry :  $E_8$     can we gauge a subgroup  $SO(9)$  ??

## 2000 : trying to construct d=2, SO(9) supergravity

- introducing vector fields ? (non-propagating in d=2)

$$\delta A_\mu^M = 2 \mathcal{V}^M{}_{IJ} (\bar{\epsilon}^I \psi_\mu^J) - i \Gamma_{A\dot{A}}^I \mathcal{V}^M{}_A (\bar{\epsilon}^I \gamma_\mu \chi^{\dot{A}})$$

we can introduce 248 (= dim E<sub>8</sub>) vector fields on which supersymmetry closes provided that  $F_{\mu\nu}^M = 0$  ! (origin : d=3)

- gauging requires Yukawa couplings and fermion shifts

$$\delta \psi_\mu^I = D_\mu \epsilon^I - g A^{IJ} \gamma_\mu \epsilon^J \quad \text{etc.}$$

with scalar tensors  $A^{IJ}(\mathcal{V}, Y_M)$  depending on scalars and dual potentials !

$$D_\mu Y_M \equiv \varepsilon_{\mu\nu} J_M^\nu$$

non-abelian duality relation

implies non-trivial consistency relations  $[F_{\mu\nu}, Y] = \varepsilon_{\mu\nu} (D_\rho J^\rho) \propto \varepsilon_{\mu\nu} \partial_\nu V_{\text{pot}}$

surprise n° 1 : consistent supersymmetric system of field equations !

# 2000 : trying to construct d=2, SO(9) supergravity

- minimal couplings and the embedding tensor

$$D_\mu \equiv \partial_\mu - g A_\mu^M \Theta_{MN} t^N$$

surprise n° 2 : group theory determines all possible gaugings !

$$\Theta_{MN} \subset \mathbf{1} + \mathbf{3875}$$

$$\text{SO}(8) \times \text{SO}(8)$$

distinguished compact gauge group

$$\text{SO}(7, 1) \times \text{SO}(7, 1)$$

$$\text{SO}(6, 2) \times \text{SO}(6, 2)$$

$$\text{SO}(5, 3) \times \text{SO}(5, 3)$$

$$\text{SO}(4, 4) \times \text{SO}(4, 4)$$

$$G_{2(2)} \times F_{4(4)}$$

$$G_2 \times F_{4(-20)}$$

$$E_{6(6)} \times \text{SL}(3)$$

$$E_{6(2)} \times \text{SU}(2, 1)$$

$$E_{6(-14)} \times \text{SU}(3)$$

$$E_{7(7)} \times \text{SL}(2)$$

$$E_{7(-5)} \times \text{SU}(2)$$

$$E_{8(8)}$$

several maximal subgroups of  $E_8$

no trace of  $\text{SO}(9)$  ... (no  $\text{SO}(9)$  singlet in  $\mathbf{3875}$ )



# 2000 : trying to construct d=2, SO(9) supergravity

- minimal couplings and the embedding tensor

$$D_\mu \equiv \partial_\mu - g A_\mu^M \Theta_{MN} t^N$$

surprise n° 2 : group theory determines all possible gaugings !

$$\Theta_{MN} \subset \mathbf{1} + \mathbf{3875}$$

SO(8) × SO(8) distinguished compact gauge group

surprise n° 3 : in fact, this is a d=3 theory !

$$D_\mu Y_M \equiv \varepsilon_{\mu\nu} J_M^\nu$$



$$\Theta_{MN} F_{\mu\nu}^N \equiv \varepsilon_{\mu\nu\rho} J_M^\rho$$

$$[F_{\mu\nu}, Y] = \varepsilon_{\mu\nu} (D_\rho J^\rho) \propto \varepsilon_{\mu\nu} \partial_\nu V_{\text{pot}}$$

non-abelian scalar-vector duality in d=3

surprise n° 4 : there is a 'simple' Lagrangian description !

$$\mathcal{L} = \text{tr} [P_\mu P^\mu] + \Theta_{MN} A^M \wedge dA^N + \dots$$

gauged sigma-model with Chern-Simons coupled vector fields !

## 2000 : trying to construct d=2, SO(9) supergravity

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$$D_\mu \equiv \partial_\mu - g A_\mu^M \Theta_{MN} t^N \quad \Theta_{MN} \subset \mathbf{1} + \mathbf{3875}$$

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SO(8)  $\times$  SO(8) distinguished compact gauge group

### d=3 maximal gauged supergravity

- construction of three-dimensional AdS supergravities
- new AdS vacua
- holography & fluxes
- construction/systematics of higher-dimensional gaugings
- tensor hierarchies of non-abelian p-forms
- rigid supersymmetry, flat target spaces

## 2000 : trying to construct d=2, SO(9) supergravity

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$$D_\mu \equiv \partial_\mu - g A_\mu^M \Theta_{MN} t^N \quad \Theta_{MN} \subset \mathbf{1} + \mathbf{3875}$$

$$\mathcal{L} = \text{tr} [P_\mu P^\mu] + \Theta_{MN} A^M \wedge dA^N + \dots$$

gauged sigma-model with Chern-Simons coupled vector fields !

SO(8)  $\times$  SO(8) distinguished compact gauge group

**d=3 maximal gauged supergravity**

where is the SO(9) theory ?

**12 years later ...**

**2012 : Golm, Hermann's 60<sup>th</sup> birthday ...**

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# 2012 : constructing $SO(9)$ supergravity

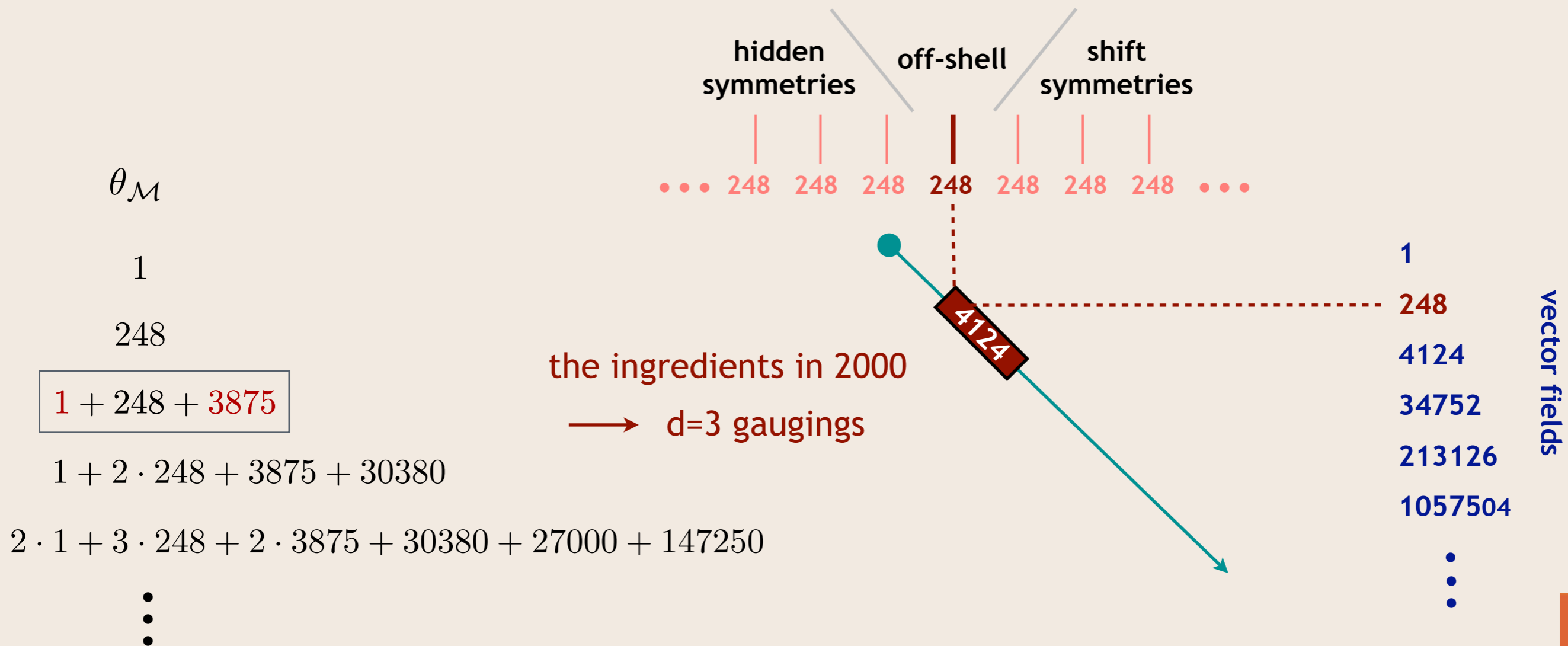
[with Thomas Ortiz]

# 2007 : gauging d=2 supergravity

$$D_\mu \equiv \partial_\mu - g A_\mu{}^{\mathcal{M}} \Theta_{\mathcal{M}}{}^A T_A$$

$$\Theta_{\mathcal{M}}{}^A = (T_B)_{\mathcal{M}}{}^{\mathcal{N}} \eta^{AB} \theta_{\mathcal{N}}$$

- d=2 supergravity has an affine symmetry group :  $E_9 = \widehat{E}_8$
- vector fields transform in the basic representation of  $E_9$
- the embedding tensor transforms in the basic representation of  $E_9$

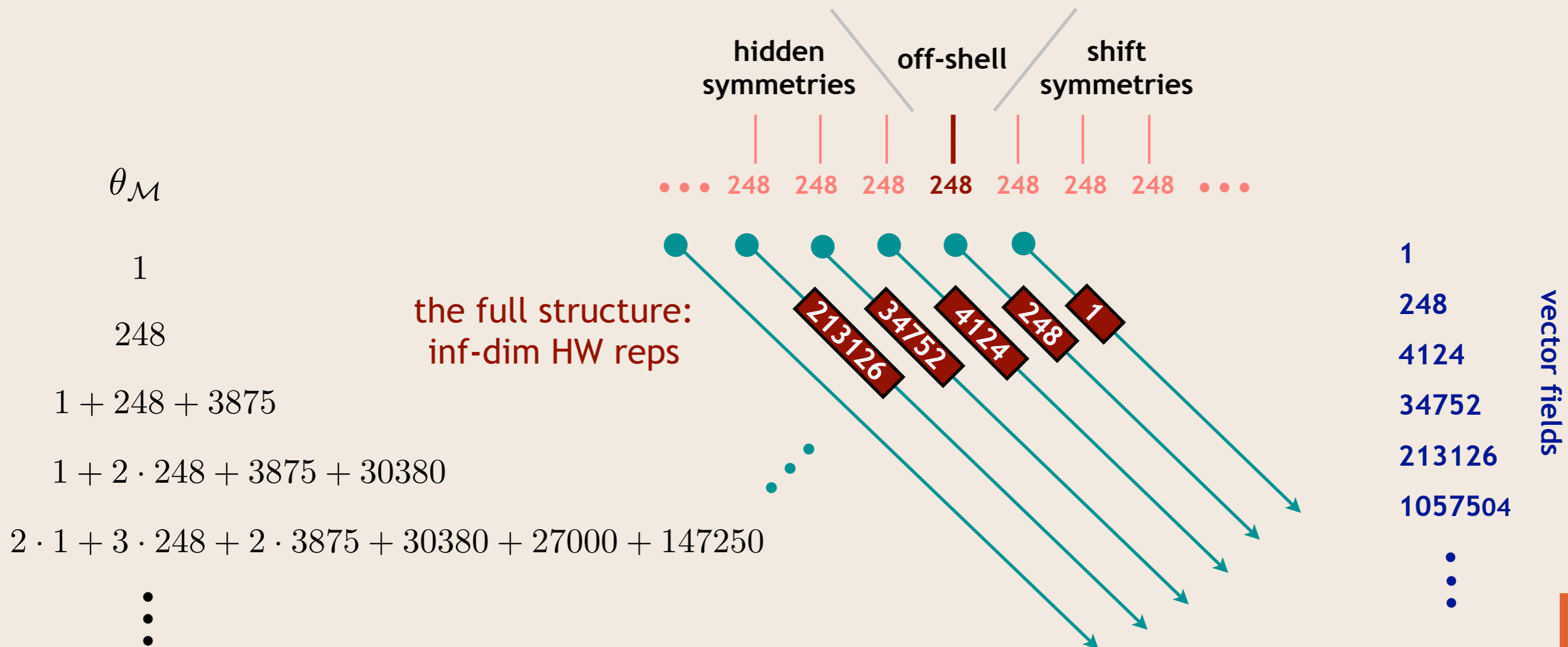


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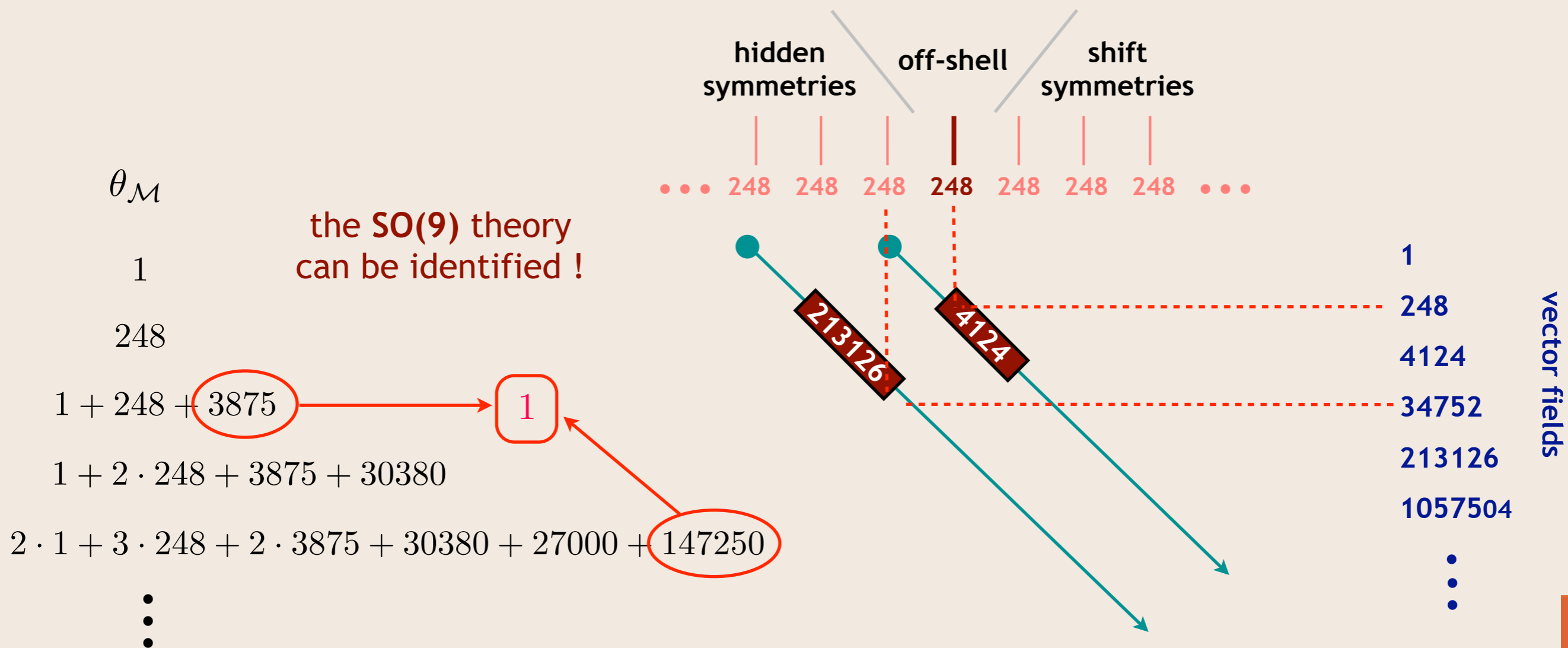


# 2012 : constructing SO(9) supergravity

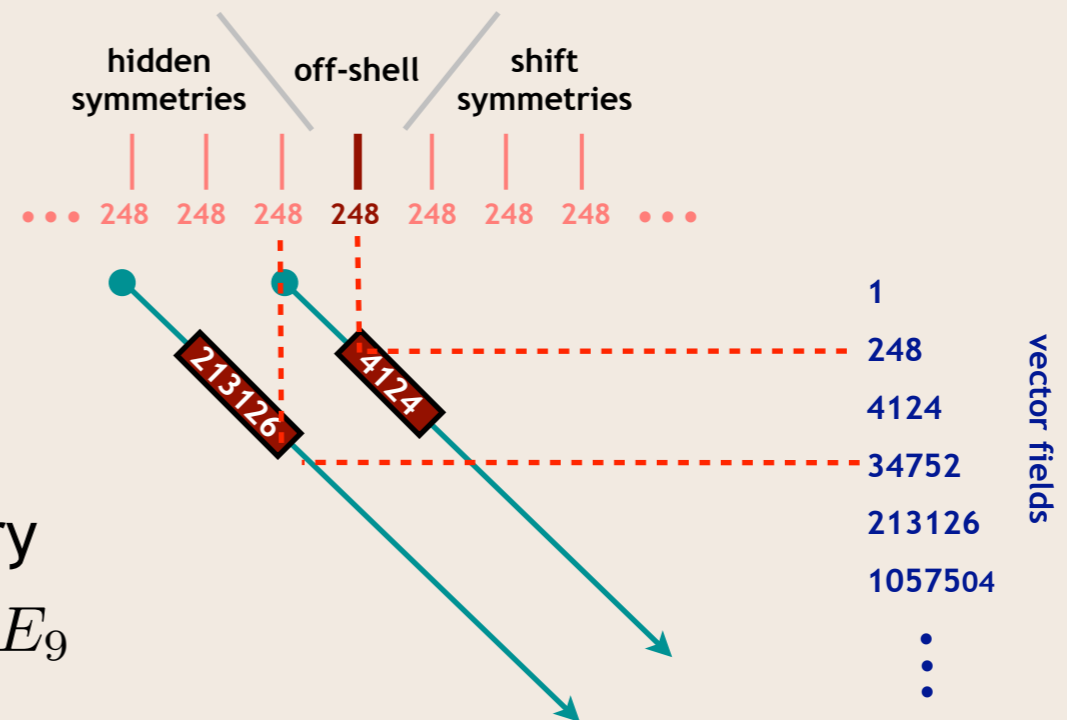
$$D_\mu \equiv \partial_\mu - g A_\mu{}^M \Theta_M{}^A T_A$$

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# 2012 : constructing SO(9) supergravity



- the SO(9) theory is a genuine d=2 theory  
in particular  $SO(9) \not\subset E_8$  but  $SO(9) \subset E_9$

- the full gauge group is infinite-dimensional (shift symmetries)

- the theory in the “E<sub>8</sub> frame” looks rather miserable

in particular the gauge group is

$$G = \underbrace{SO(8) \times \left( (\mathbb{R}_+^{28} \times \mathbb{R}_+^8)_0 \right)}_{\text{off-shell}} \times \underbrace{(\mathbb{R}_+^8)_{-1}}_{\text{hidden (on-shell)}}$$

- go to a “T-dual frame” in which SO(9) is among the off-shell symmetries



# 2012 : constructing SO(9) supergravity

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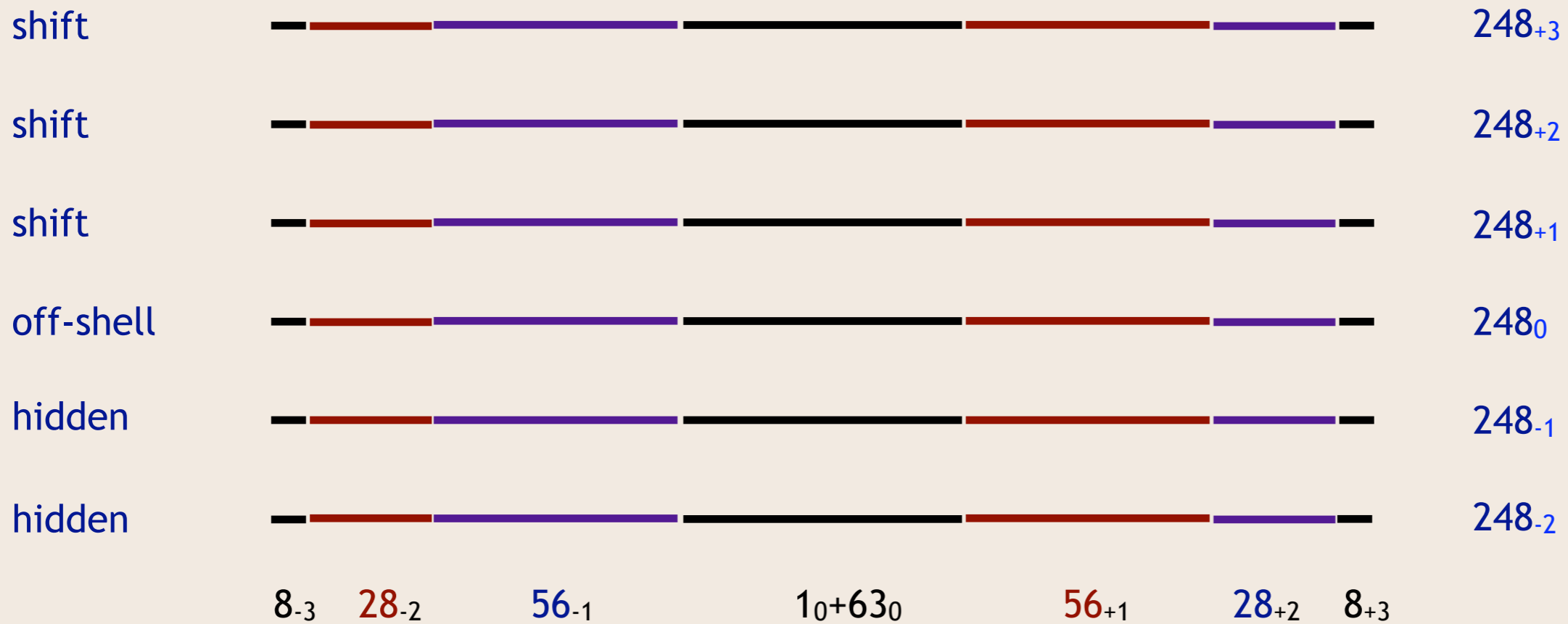
affine  $E_8$  with  $L_0$  grading

shift	—————	$248_{+3}$
shift	—————	$248_{+2}$
shift	—————	$248_{+1}$
off-shell	—————	$248_0$
hidden	—————	$248_{-1}$
hidden	—————	$248_{-2}$

# 2012 : constructing SO(9) supergravity

affine  $E_8$  with  $L_0$  grading

breaking under  $SL(8) \times \mathbb{R}^+$



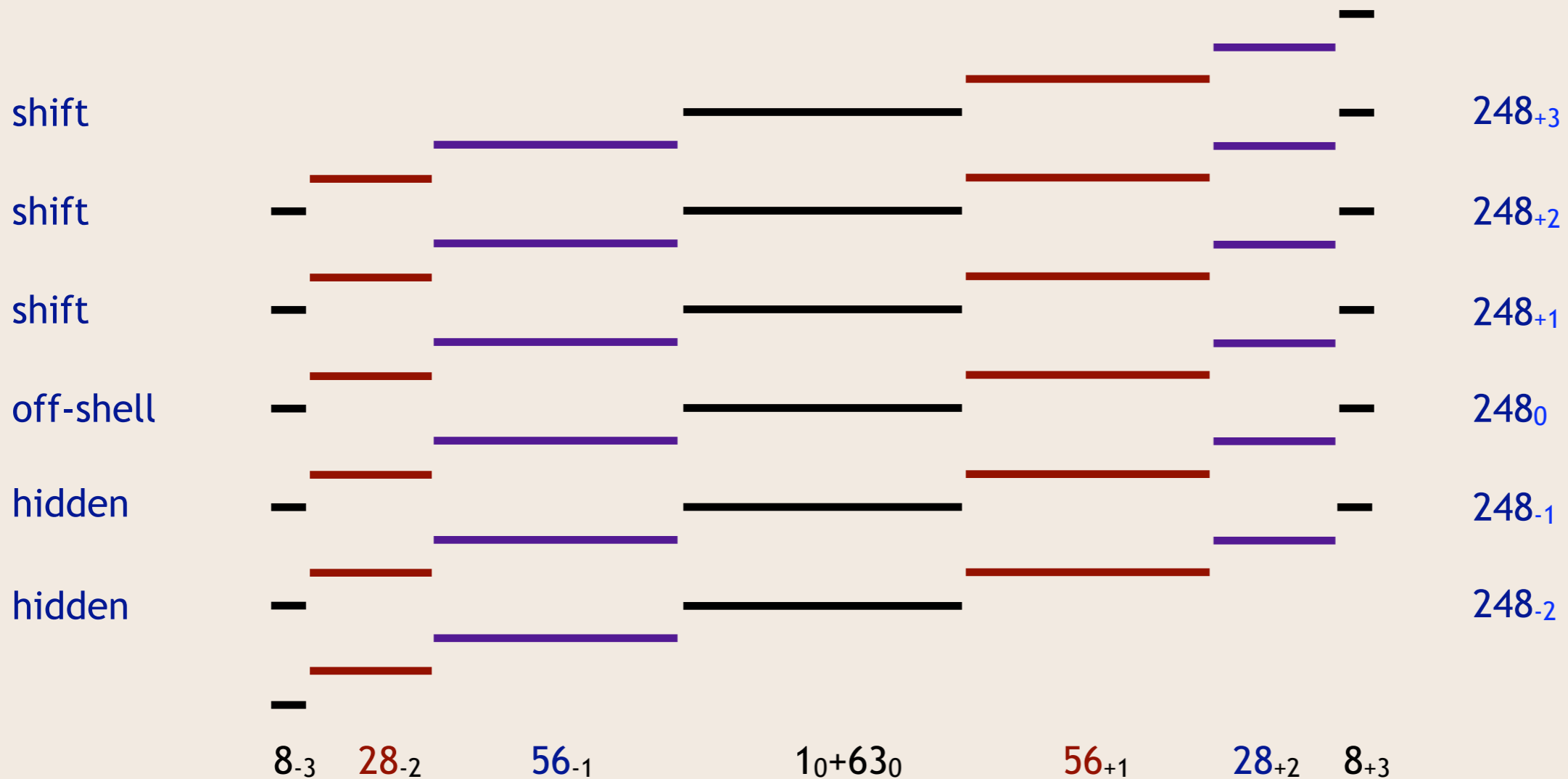
diagonal grading :  $\widehat{L}_0 \equiv L_0 + \frac{1}{3}q_{\mathbb{R}}$



# 2012 : constructing SO(9) supergravity

affine  $E_8$  with  $L_0$  grading

breaking under  $SL(8) \times \mathbb{R}^+$

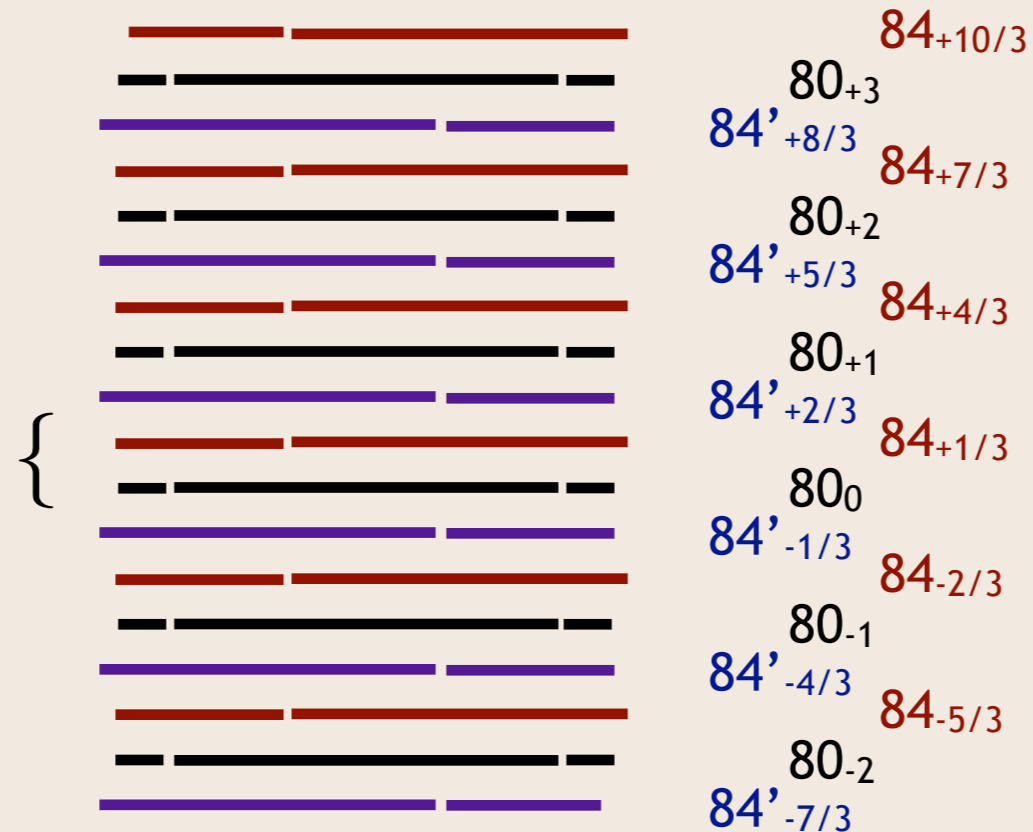


diagonal grading :  $\widehat{L}_0 \equiv L_0 + \frac{1}{3}q_{\mathbb{R}}$

# 2012 : constructing SO(9) supergravity

affine  $E_8$  with  $\widehat{L}_0$  grading : decomposition under  $\widehat{SL}(9)$

off-shell symmetry  
 $SL(9) \ltimes \mathbb{T}^{84}$



“T-dual frame” :

coset sigma model  $E_8/SO(16)$   $\longrightarrow$  coset sigma model  $(SL(9) \ltimes \mathbb{T}^{84})/SO(9)$   
 with WZW term

# 2012 : constructing SO(9) supergravity

“T-dual frame” :

coset sigma model  $(SL(9) \ltimes \mathbb{T}^{84}) / SO(9)$  with WZW term  $84 \wedge 84 \wedge 84 \longrightarrow 1$

$$\begin{aligned} \mathcal{L}_0 = & -\frac{1}{4}\rho R + \frac{1}{4}\rho P^{\mu ab} P_{\mu}^{ab} + \frac{1}{12}\rho^{1/3} M_{il} M_{jm} M_{kn} \partial^{\mu} \phi^{ijk} \partial_{\mu} \phi^{lmn} \\ & + \frac{1}{648} \varepsilon^{\mu\nu} \varepsilon_{klmnpqrst} \phi^{klm} \partial_{\mu} \phi^{npq} \partial_{\nu} \phi^{rst} \end{aligned}$$

in fact this is the d=11 theory reduced on a torus  $T^9$  ...

fermionic part :

$$\begin{aligned} & -\rho e^{-1} \varepsilon^{\mu\nu} \bar{\psi}_2^I D_{\mu} \psi_{\nu}^I - \frac{i}{2} \bar{\psi}_{\nu}^I \gamma^{\nu} \psi_{\mu}^I \partial^{\mu} \rho - \frac{i}{2} \rho \bar{\chi}^{aI} \gamma^{\mu} D_{\mu} \chi^{aI} + \frac{i}{2} \rho^{2/3} \bar{\chi}^{aI} \gamma^3 \gamma^{\mu} \chi^{bJ} \Gamma_{IJ}^c \varphi_{\mu}^{abc} - \frac{i}{24} \rho^{2/3} \bar{\chi}^{aI} \gamma^3 \gamma^{\mu} \chi^{aJ} \Gamma_{IJ}^{bcd} \varphi_{\mu}^{bcd} \\ & - \frac{1}{4} \rho^{2/3} \bar{\chi}^{aI} \gamma^3 \gamma^{\nu} \gamma^{\mu} \psi_{\nu}^J \Gamma_{IJ}^{bc} \varphi_{\mu}^{abc} - \frac{i}{12} \rho^{2/3} \bar{\chi}^{aI} \gamma^{\mu} \psi_2^J \Gamma_{IJ}^{bc} \varphi_{\mu}^{abc} - \frac{1}{2} \rho \bar{\chi}^{aI} \gamma^{\nu} \gamma^{\mu} \psi_{\nu}^J \Gamma_{IJ}^b P_{\mu}^{ab} - \frac{i}{2} \rho \bar{\chi}^{aI} \gamma^3 \gamma^{\mu} \psi_2^J \Gamma_{IJ}^b P_{\mu}^{ab} \\ & + \frac{i}{54} \rho^{2/3} \bar{\psi}_2^I \gamma^3 \gamma^{\mu} \psi_2^J \Gamma_{IJ}^{abc} \varphi_{\mu}^{abc} + \frac{1}{24} \rho^{2/3} \bar{\psi}_2^I \left( \gamma^{\mu} \gamma^{\nu} - \frac{1}{3} \gamma^{\nu} \gamma^{\mu} \right) \psi_{\nu}^J \Gamma_{IJ}^{abc} \varphi_{\mu}^{abc} \end{aligned}$$

off-shell symmetry  $SL(9) \ltimes \mathbb{T}^{84}$

$\supset SO(9)$  gauging

# 2012 : constructing SO(9) supergravity

“T-dual frame” :

**gauged** coset sigma model  $(SL(9) \times \mathbb{T}^{84}) / SO(9)$  with WZW term

$$\mathcal{L} = -\frac{1}{4}\rho R + \frac{1}{4}\rho P^{\mu ab} P_{\mu}^{ab} + \frac{1}{12}\rho^{1/3} M_{il} M_{jm} M_{kn} D^{\mu} \phi^{ijk} D_{\mu} \phi^{lmn} + \frac{1}{648} \varepsilon^{\mu\nu} \varepsilon_{klmnpqrst} \phi^{klm} D_{\mu} \phi^{npq} D_{\nu} \phi^{rst}$$

fermion couplings and Yukawa terms

$$\begin{aligned} \mathcal{L}_{\text{Yuk}} &= -\frac{1}{2} e^{-1} \rho \varepsilon^{\mu\nu} \left( \bar{\psi}_{\nu}^I \psi_{\mu}^J B_{IJ} + \bar{\psi}_{\nu}^I \gamma^3 \psi_{\mu}^J \tilde{B}_{IJ} - 2i \bar{\psi}_2^I \gamma_{\nu} \psi_{\mu}^J A_{IJ} \right) + i \rho \bar{\psi}_2^I \gamma^{\mu} \psi_{\mu}^J \tilde{A}_{IJ} \\ &\quad + i \rho \bar{\chi}^{aI} \gamma^{\mu} \psi_{\mu}^J C_{IJ}^a - i \rho \bar{\chi}^{aI} \gamma^3 \gamma^{\mu} \psi_{\mu}^J \tilde{C}_{IJ}^a + \rho \bar{\psi}_2^I \psi_2^J D_{IJ} + \rho \bar{\psi}_2^I \gamma^3 \psi_2^J \tilde{D}_{IJ} \\ A_{IJ} &= \frac{7}{9} \delta_{IJ} b - \frac{1}{9} \Gamma_{IJ}^a b^a + \frac{1}{9} \Gamma_{IJ}^{abcd} b^{abcd}, \quad b = \frac{1}{4} \rho^{-2/9} T, \\ \tilde{A}_{IJ} &= \frac{2}{9} \Gamma_{IJ}^{ab} b^{ab} - \frac{4}{9} \Gamma_{IJ}^{abc} b^{abc}, \quad b^a = -\rho^{-14/9} \mathcal{V}^{-1km}{}_{bc} \theta_{ml} \varphi^{abc} Y_k^l + \frac{1}{144} \rho^{-14/9} \varepsilon^{bcdefghij} T^{kl} \varphi^{kef} \varphi^{lgh} \varphi^{aij} \varphi^{bcd}, \\ B_{IJ} &= \Gamma_{IJ}^{ab} b^{ab} + \Gamma_{IJ}^{abc} b^{abc}, \quad \Gamma_{IJ}^{ab} = \frac{1}{9} \rho^{-11/9} \mathcal{V}^{-1[km]}{}_{ab} \theta_{ml} Y_k^l + \frac{1}{144} \rho^{-11/9} \varepsilon^{bcdefghij} T^{jk} \varphi^{jcd} \varphi^{kef} \varphi^{ghi}, \\ \tilde{B}_{IJ} &= \delta_{IJ} b + \Gamma_{IJ}^a b^a + \Gamma_{IJ}^{abcd} b^{abcd}, \quad \Gamma_{IJ}^{abcd} = \frac{1}{9} \rho^{-5/9} \Gamma_{IJ}^{abcd} \varphi^{bcd}, \\ C_{IJ}^a &= \frac{8}{9} \delta_{IJ} b^a - \frac{1}{9} \Gamma_{IJ}^{ab} b^b + \frac{20}{9} \Gamma_{IJ}^{abcd} b^{abcd} - \frac{1}{9} \Gamma_{IJ}^{abcde} b^{abcde} + \frac{1}{9} \Gamma_{IJ}^{abcde} b^{abcde} - \frac{1}{9} \Gamma_{IJ}^{abcde} b^{abcde} - 12 \Gamma_{IJ}^{cd} b^{abcd} - 2 c^{ab} \delta_{IJ} \\ \tilde{C}_{IJ}^a &= -\frac{14}{9} \Gamma_{IJ}^b b^{ab} + \frac{2}{9} \Gamma_{IJ}^{abc} b^{abc} + \frac{2}{9} \Gamma_{IJ}^{bc} b^{bc} - \frac{1}{9} \Gamma_{IJ}^{abcd} b^{abcd} + \frac{1}{9} \Gamma_{IJ}^{abc} b^{abc} - \frac{1}{9} \Gamma_{IJ}^{abc} b^{abc} - \frac{1}{9} \Gamma_{IJ}^{abc} b^{abc} - 2 c^{c,ab} \Gamma_{IJ}^c, \\ D_{IJ} &= \frac{14}{81} \delta_{IJ} b - \frac{70}{81} \Gamma_{IJ}^a b^a + \frac{28}{81} \Gamma_{IJ}^{abcd} b^{abcd}, \quad c^{a,bc} = \frac{1}{3} \rho^{-5/9} (T^{da} \varphi^{bcd} + T^{d[b} \varphi^{c]da}), \end{aligned}$$

# 2012 : constructing SO(9) supergravity

“T-dual frame” :

**gauged** coset sigma model  $(SL(9) \ltimes \mathbb{T}^{84}) / SO(9)$  with WZW term

$$\mathcal{L} = -\frac{1}{4}\rho R + \frac{1}{4}\rho P^{\mu ab} P_{\mu}^{ab} + \frac{1}{12}\rho^{1/3} M_{il} M_{jm} M_{kn} D^{\mu} \phi^{ijk} D_{\mu} \phi^{lmn} \\ + \frac{1}{648} \varepsilon^{\mu\nu} \varepsilon_{klmnpqrst} \phi^{klm} D_{\mu} \phi^{npq} D_{\nu} \phi^{rst}$$

vector fields couple via

$$\mathcal{L}_F = \varepsilon^{\mu\nu} F_{\mu\nu}{}^{mn} \mathcal{Y}_{mn} \quad \text{with auxiliary (dual scalar) fields } \mathcal{Y}_{mn}$$

scalar potential

$$V_{\text{pot}} = \frac{1}{8} \rho^{5/9} \left( (\text{tr } T)^2 - 2 \text{tr}(T^2) + 18 \rho^{-2/3} T^{d[a} \varphi^{bc]d} T^{ea} \varphi^{bce} - 16 \rho^{-2/3} T^{d[b} \varphi^{c]ad} T^{eb} \varphi^{cae} \right)$$

$$- \rho^{-13/9} T^{ac} T^{bc} Y_{ad} Y_{bd} + \mathcal{O}(\phi^3)$$

eighth order polynomial in  $\phi$

$$T \equiv (\mathcal{V}^T \mathcal{V})^{-1}$$

$$\varphi \equiv \phi \cdot \mathcal{V}$$

$$Y \equiv \mathcal{V}^T \mathcal{Y} \mathcal{V}$$

the dilaton powers precisely support the correct DW solution (near horizon of  $AdS_2 \times S^8$ )

# 2012 : constructing SO(9) supergravity

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different presentations

**gauged** coset sigma model  $(SL(9) \ltimes \mathbb{T}^{84}) / SO(9)$  with WZW term

$$\begin{aligned} \mathcal{L} = & -\frac{1}{4}\rho R + \frac{1}{4}\rho P^{\mu ab} P_{\mu}^{ab} + \frac{1}{12}\rho^{1/3} M_{il} M_{jm} M_{kn} D^{\mu} \phi^{ijk} D_{\mu} \phi^{lmn} \\ & + \frac{1}{648} \varepsilon^{\mu\nu} \varepsilon_{klmnpqrst} \phi^{klm} D_{\mu} \phi^{npq} D_{\nu} \phi^{rst} + \varepsilon^{\mu\nu} F_{\mu\nu}{}^{mn} \mathcal{Y}_{mn} + V_{\text{pot}}(\rho, \mathcal{V}, \phi, \mathcal{Y}) \end{aligned}$$

---

integrate out the vector fields  $A_{\mu}{}^{mn}$

$$\mathcal{L}_{\text{T}} = -\frac{1}{4}\rho R + G_{ij}(\rho, \mathcal{V}, \phi, \mathcal{Y}) \partial_{\mu} \Phi^i \partial^{\mu} \Phi^j + \varepsilon^{\mu\nu} B_{ij}(\rho, \mathcal{V}, \phi, \mathcal{Y}) \partial_{\mu} \Phi^i \partial_{\nu} \Phi^j + V_{\text{pot}}(\rho, \mathcal{V}, \phi, \mathcal{Y})$$

ungauged sigma model on a different target space

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integrate out the auxiliary scalars  $\mathcal{Y}_{mn}$

$$\mathcal{L}_2 = -\frac{1}{4}\rho R + F_{\mu\nu}{}^{mn} F^{\mu\nu}{}^{kl} \mathcal{R}_{mn,kl}(\rho, \mathcal{V}, \phi) + \dots + \tilde{V}_{\text{pot}}(\rho, \mathcal{V}, \phi)$$

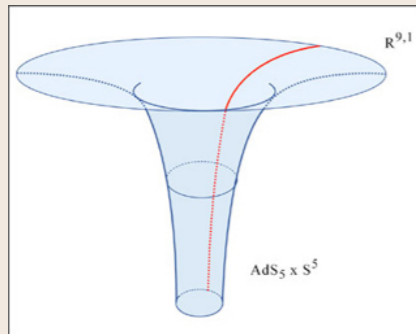
gauged sigma model coupled to d=2 SYM



# concluding

## SO(9) supergravity

- maximally supersymmetric d=2 supergravity with gauge group SO(9)
- last missing gauged supergravity around near-horizon geometries



warped

D6	IIA	$AdS_8 \times S^2$	d=8, SO(3)	[Salam, Sezgin, 1984]
D5	IIB	$AdS_7 \times S^3$	d=7, SO(4)	[Samtleben, Weidner, 2005]
D4	IIA	$AdS_6 \times S^4$	d=6, SO(5)	[Pernici, Pilch, van Nieuwenhuizen, 1984]
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D2	IIA	$AdS_4 \times S^6$	d=4, SO(7)	[Hull, 1984]
F1/D1	IIA/B	$AdS_3 \times S^7$	d=3, SO(8)	[de Wit, Nicolai, 1982]
D0	IIA	$AdS_2 \times S^8$	d=2, SO(9)	

- mission completed ...
- holography : d=1 supersymmetric matrix quantum mechanics ...!

# concluding

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## SO(9) supergravity

- maximally supersymmetric d=2 supergravity with gauge group SO(9)
- **last missing gauged supergravity around near-horizon geometries**
- holography : d=1 supersymmetric matrix quantum mechanics ...!

we came here by a wonderful detour via three (and many other) dimensions

with Hermann via the

**d=3, SO(8) x SO(8) theory**

which is still awaiting its embedding/interpretation in higher dimensions,  
string theory, holography, ... (matrix string theory, double field theory, ...???)

**sufficient material for the next anniversaries ...**

**Happy Birthday, Hermann ! (in all dimensions)**