

Random tensors

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Orsay, Quantum gravity in Paris

Introduction to random tensors

- ▶ Tensor $T_{a_1 \dots a_d}$ is an array of complex numbers of size $N \times N \times \dots \times N = N^d$
- ▶ **Goal:** Find the **typical** properties of tensors at large N .

Example

Wigner and the spectrum of heavy nuclei

Consider Hamiltonian as random Hermitian matrix

What is the eigenvalue density?

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Example

Disordered systems, like spin glasses

Spins S_a on lattice sites a

Take d spins interacting together, $H = - \sum_{\text{sites}} T_{a_1 \dots a_d} S_{a_1} \dots S_{a_d}$

VB, Gurau, Smerlak, arXiv:1206.5539

Summing over lattices

- ▶ Family of 'lattices' (e.g. triangulations) $\{G\}$, with amplitude $A(G)$

$$\text{Generating function} \quad \sum_{\{G\}} \lambda^{\#\text{vertices}(G)} A(G)$$

- ▶ Interpret this as a Feynman expansion of some integral

2d quantum gravity

Sum over quadrangulated surfaces = integral over (hermitian) matrices

$$\int [dM_{N \times N}] e^{-N(\text{tr} M^2 + \lambda \text{tr} M^4)} \underset{\text{large } N}{=} \sum_{\substack{\text{quadrangulations} \\ \text{of the sphere}}} (-\lambda)^{V(G)} \frac{1}{s(G)}$$

- ▶ Invariant quantities $\text{tr} M^p$: loop observables

$$\langle \text{tr} M^p \rangle = \text{disc amplitude with } p \text{ boundary segments}$$

- ▶ Exists for more generic potentials than M^4

Matrix models

A combinatorial tool, surprisingly powerful

- ▶ Add matter degrees of freedom to the amplitude
- ▶ Generating functions of discretized surfaces with matter
- ▶ Exactly solve enumerative problem in stat. mech.
- ▶ Get 2d quantum gravity coupled to matter (= non-critical strings)
Developed by David, Kazakov and many others in the 80s-90s
- ▶ Related to conformal field theories in several ways
- ▶ Schwinger-Dyson equations exhibit a Virasoro algebra
- ▶ Get integrable hierarchies (Volterra, KdV)

Methods

- ▶ Integral over eigenvalues
- ▶ Schwinger-Dyson equations
- ▶ Pure combinatorics
Modern developments with Schaeffer's bijection, integrable combinatorics

Tensors generalize matrices

The historical way: random triangulations of manifolds

- ▶ Sum triangulations in d dimensions using **rank d tensors**
Generalize matrix M_{ab} to T_{a_1, \dots, a_d}
Suggested by Ambjorn, Gross, Sasakura in '91 independently.
- ▶ But no way to get analytical results, until 2011 !
- ▶ Numerical simulations (fixed topology): Euclidean Dynamical Triangulations [**and better: Renate Loll's talk**]
- ▶ In group field theories (Boulatov, Ooguri, '90s), one sums over cell complexes too, with different amplitudes
[**Daniele Oriti's talk**]
- ▶ Control sum over triangulations via Feynman rules only
- ▶ No eigenvalues
- ▶ Solely rely on **Schwinger-Dyson eqs** and **combinatorics**

Objective I

- ▶ Extract an **exactly solvable** sector of dynamical triangulations
- ▶ That works in arbitrary dimensions

Reproduce the basic results of matrix models

- ▶ Show that the large N limit exists
- ▶ Show that it projects onto a summable family of triangulations
The **melon** family
- ▶ Extract singularities of the generating functions
Continuum limit and multicritical behaviors
- ▶ Prove universality of continuum limits
 - ▶ irrespective of the microscopic details
(triangulations, quadrangulations, etc.)
 - ▶ find how to change the the microscopic details from tensors!

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Program started in 2011, **Done**

References

- ▶ Critical behavior of colored tensor models in the large N limit
VB, R. Gurau, A. Riello, V. Rivasseau, Nucl. Phys. B (2011)
- ▶ The Ising Model on Random Lattices in Arbitrary Dimensions
VB, R. Gurau, V. Rivasseau, Phys. Lett. B (2011)
- ▶ Multicritical tensor models and hard dimers on spherical random lattices, *VB*, Phys. Lett. A (2012)
- ▶ Random tensor models in the large N limit: Uncoloring the colored tensor models, *VB, R. Gurau, V. Rivasseau*, Phys. Rev. D (2012)
- ▶ Hard dimers coupled to random tensors, *VB, H. Erbin*, JSTAT (2012)
- ▶ Counting Line-Colored D-ary Trees, *VB, R. Gurau*, (2012)
- ▶ Phase Transition in Dually Weighted Colored Tensor Models
R. Gurau, D. Benedetti, Nucl. Phys. B (2012)
- ▶ Melons are branched polymers
R. Gurau, J. Ryan, (2013)

Objective II

- ▶ Melonic geometry: random trees [Come to J. Ryan's talk!]

Further explore the richness of tensor models

- ▶ More invariant objects = boundary triangulations, than in matrix models!
- ▶ Landscape of large N behaviors
New $1/N$ expansions in random tensor models, VB, arXiv:1211.1657
- ▶ Universality theorem vs non-universal models
- ▶ Study the Schwinger-Dyson equations
Revisiting random tensor models at large N via the Schwinger-Dyson equations, VB, JHEP
- ▶ Complicated set of equations labeled by boundary triangulations
- ▶ Relations to combinatorics
(meanders aka polymer folding [to appear])
- ▶ Form an infinite-dimensional Lie algebra!
Reduces to Virasoro in the matrix case
Conformal aspects, integrability?

Colored graphs and unitary-invariant tensor ensemble

Colored Graphs and Triangulations

Tensor models and their large N expansion

Schwinger-Dyson equations and Universality

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Generating graphs with integrals

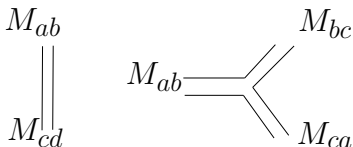
$$Z(\lambda, \alpha) = \sum_{\text{Graphs } \{G\}} \lambda^V Z_{G_n}(\alpha)$$

Interpret it as the Feynman expansion of some integrals

Choice of variables depends on the family of graphs under consideration.

Example: Matrix integrals generate ribbon graphs

$$[dM] \underbrace{e^{-N \text{tr } M^2}}_{\text{Gaussian}} \underbrace{e^{-Ng \text{tr } M^4}}_{\text{perturbation}}$$



Stranded graphs and tensors

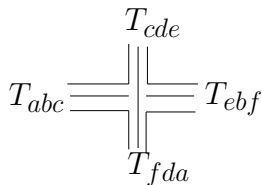
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Interpret it as the Feynman expansion of some integrals

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Example: Tensor integrals generate stranded graphs

Propagator of a Gaussian measure
Interaction



But no way to control these models analytically!

Tensor Ensemble

External tensor product of $U(N)$ acts like

$$T'_{a_1 \dots a_d} = \sum_{b_1, \dots, b_d} U_{a_1 b_1}^{(1)} \dots U_{a_d b_d}^{(d)} T_{b_1 \dots b_d}$$

and on c.c. \bar{T} .

Invariant monomials

Contract index in position i of T with an index in position i in a \bar{T} .

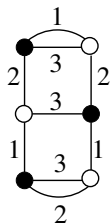
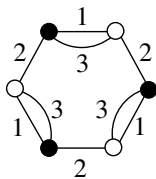
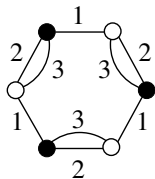
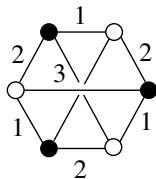
Equivalence with colored graphs

- ▶ a T → white vertex
- ▶ a \bar{T} → black vertex
- ▶ Index → half-line on a vertex
- ▶ Position of index → color label on half-line
- ▶ Index contraction → connect half-lines with same color

Colored graphs

d -colored, connected graphs

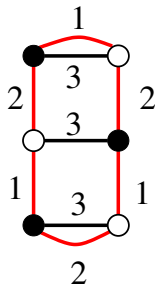
- ▶ Bipartite graphs: black and white vertices
- ▶ Colored lines with d possible colors
- ▶ d -valent nodes with distinct colors



Colored graphs

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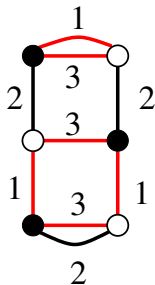


Faces are closed cycles with only **two** colors.

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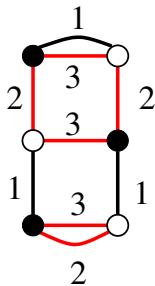


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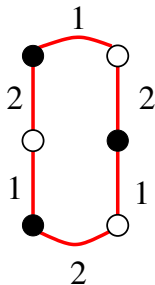


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$d = 2$ reduces to loops with 1 face and $2p$ vertices, i.e. $\text{tr}(TT^\dagger)^p$

Richer structure with tensors than with matrices

Colored graphs and unitary-invariant tensor ensemble

Colored Graphs and Triangulations

Tensor models and their large N expansion

Schwinger-Dyson equations and Universality

Relation between graphs and topology

Graphs encode topology of triangulations

- ▶ Triangulation: gluing of simplices (triangles, tetrahedra, pentachora, d -simplices), defined by attaching maps
- ▶ Δ -colored graph \leftrightarrow $(\Delta - 1)$ -triangulation

Graph dual to the triangulation

vertex	\rightarrow	$(\Delta - 1)$ -simplex
line	\rightarrow	$(\Delta - 2)$ -simplex
face	\rightarrow	$(\Delta - 3)$ -simplex
k -bubble	\rightarrow	$(\Delta - k - 1)$ -simplex

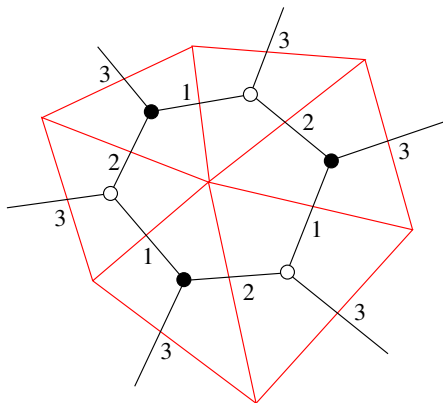
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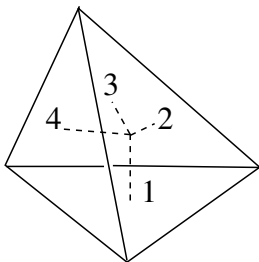
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Triangulations from colored graphs

A Δ -colored vertex is dual to a $(\Delta - 1)$ -simplex. The legs of the vertex are dual to the boundary $(\Delta - 2)$ -simplices.

- ▶ Boundary triangles labeled by a color $c = 1, \dots, \Delta$

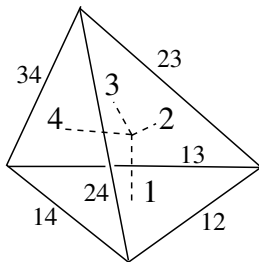


Colors identify all sub-simplices

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- ▶ Induced colorings
- ▶ Edges labeled by pair of colors

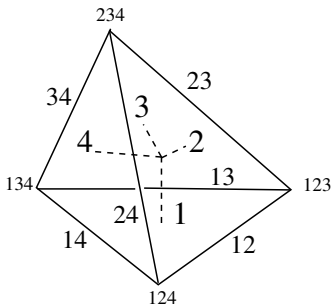


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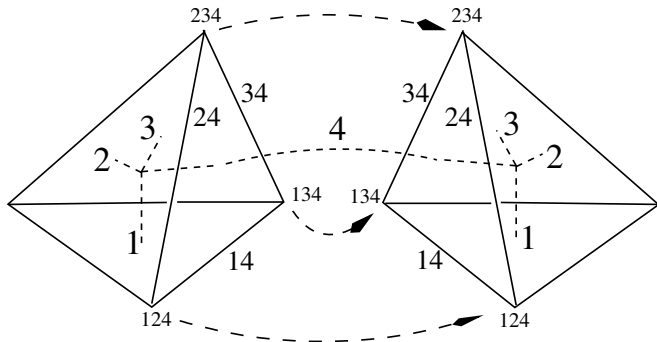
- ▶ Boundary triangles labeled by a color $c = 1, \dots, \Delta$
- ▶ Induced colorings
- ▶ Edges labeled by pair of colors
- ▶ Nodes labeled by three colors



Colors identify all sub-simplices

Triangulations from colored graphs

A propagator glues two tetrahedra along a triangle, respecting the induced coloring of all sub-simplices



Theory of crystallization and GEMs (graph-encoding manifolds):
 Δ -colored graphs are dual to triangulations of pseudo-manifolds of dimension $\Delta - 1$ [Pezzana, Ferri, Cagliardi, Lins, Caravelli].

Colored graphs and unitary-invariant tensor ensemble

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Tensor Models

- ▶ $(B_i)_{i \in I}$ finite set of bubbles. Set action to

$$S(T, \bar{T}) = T \cdot \bar{T} + \sum_{i \in I} t_i B_i(T, \bar{T})$$

- ▶ Free energy F

$$\exp -F(\{t_i\}) = Z = \int [dT d\bar{T}] \exp -\frac{N^\alpha}{\lambda} S(T, \bar{T})$$

- ▶ Bubble expectation value

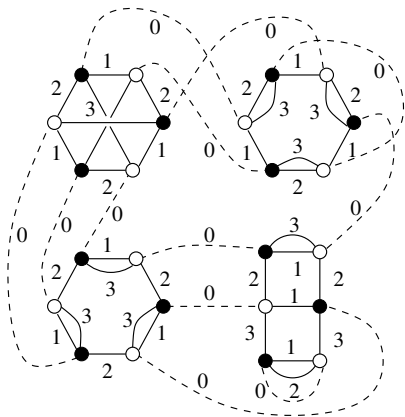
$$\langle B \rangle = \frac{1}{Z} \int [dT d\bar{T}] B(T, \bar{T}) \exp -\frac{N^\alpha}{\lambda} S(T, \bar{T})$$

- ▶ In particular $\langle B_i(T, \bar{T}) \rangle \sim \frac{\partial F}{\partial t_i}$

Feynman expansion and colored graphs

Feynman expansion of F

- ▶ Connect bubbles B_i with lines (the covariance)
- ▶ Each vertex receives such a line
- ▶ Give it color 0
- ▶ Generate $(d + 1)$ -colored, connected graphs
- ▶ Dual to d -triangulations



$$F = \sum_{(d+1)\text{-colored graphs } G} N^{\text{power}(G)} A_G(\lambda, \{t_i\})$$

1/N expansion

Connected colored graph G : f faces, l lines of color 0, b bubbles

$$A(G) \sim N^{f-\alpha(l-b)} \lambda^{l-b}$$

Value of α fixed by **balance between f and $(l-b)$** so that

- ▶ $A(G) \leq K N^{\text{fixed exponent}}$
If α too small, no bound
- ▶ If α too large, then suppression at large number of vertices, hence no infinite family, no continuum limit

Counting of faces in colored graphs

$$f - (d-1)(l-b) = d - \underbrace{\omega(G)}_{\substack{\text{positive integer} \\ \text{the degree}}} \leq d$$

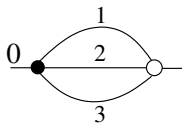
$d = 2$, the degree is the genus

Set $\alpha = d - 1$

Dominant graphs

$$f - (d - 1)(l - b) \leq d$$

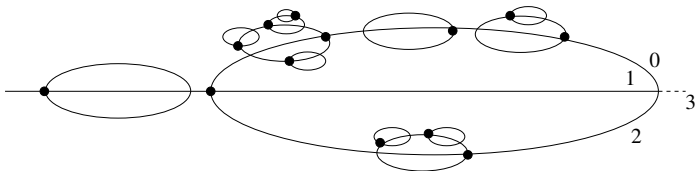
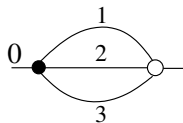
- ▶ Leading order: maximize the number of faces
- ▶ $d = 2$, get planar graphs
- ▶ $d \geq 3$, get **melon** graphs!
- ▶ Movements which maximize f
Add only two vertices, create $\frac{d(d-1)}{2}$ faces



Dominant graphs

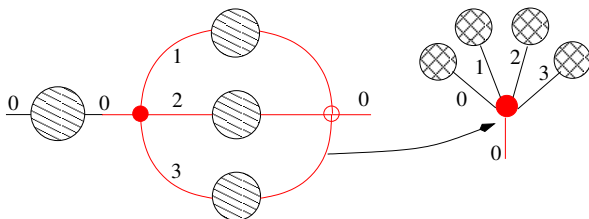
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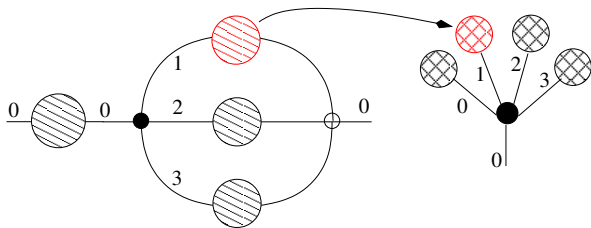
Melonic graphs have melonic bubbles
Only them are relevant in the action

Mapping melons to trees

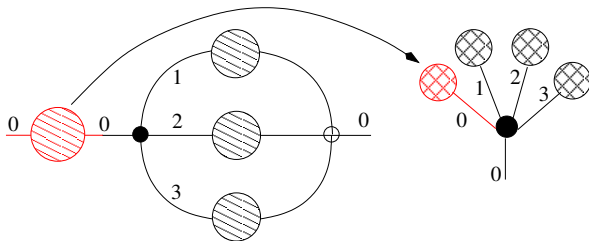


The black and white vertices of external legs are mapped to a tree vertex.
The lines connected to the black vertex are mapped to tree lines which inherit colors.

Mapping melons to trees

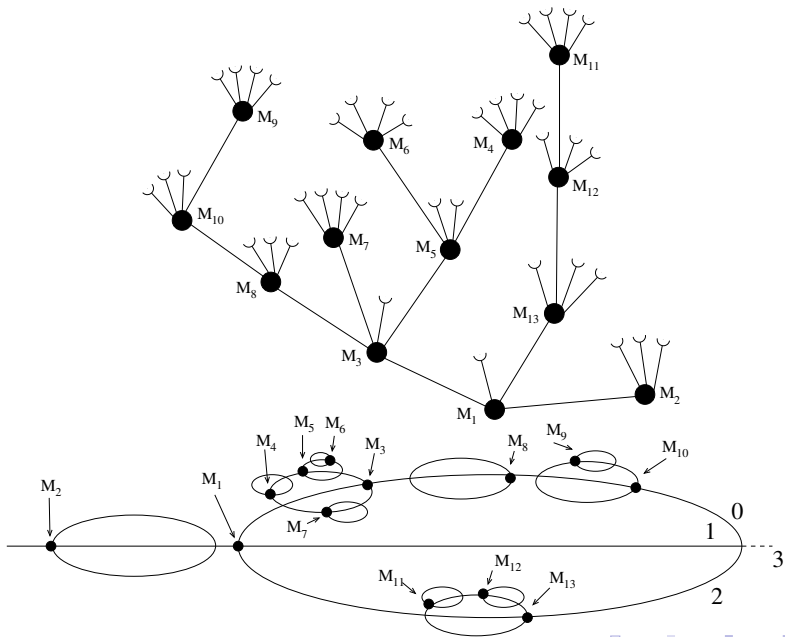


Mapping melons to trees



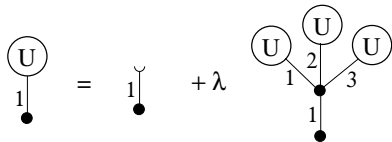
Get rooted $(d + 1)$ -ary trees, a species of Branched Polymers (BP).

Mapping melons to trees



Counting melons

- ▶ Generating functions for d -colored melonic graphs with two external lines, $U(\lambda) = \sum_p \lambda^p C_p^{(d)}$



$$U = 1 + \lambda U^d$$

$$\Rightarrow C_p^{(d)} = \frac{1}{d p + 1} \binom{d p + 1}{p}$$

Generalized Catalan numbers

- ▶ Trees with exactly p_i lines of color i

$$C_{p_1 \dots p_d}^{(d)} = \frac{1}{1 + \sum_{i=1}^d p_i} \prod_{j=1}^d \binom{1 + \sum_{i=1}^d p_i}{p_j}$$

Generalized Narayana numbers

Counting line-colored D -ary trees, VB and R. Gurau, arXiv:1206.4203
[math-ph]

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Tensor models and their large N expansion

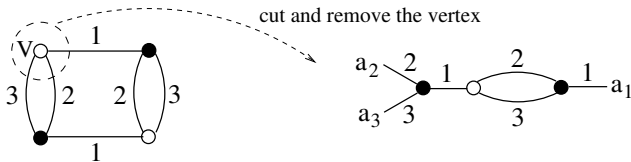
Schwinger-Dyson equations and Universality

SDEs: Quantum equations of motion

- ▶ To solve a generic model, need systematic eqs

$$\int [dTd\bar{T}] \frac{\partial}{\partial \bar{T}_{a_1 \dots a_d}} \left(\mathcal{O}(T, \bar{T}) e^{-N^{d-1} S(T, \bar{T})} \right) = 0$$

Insertions: Open bubbles

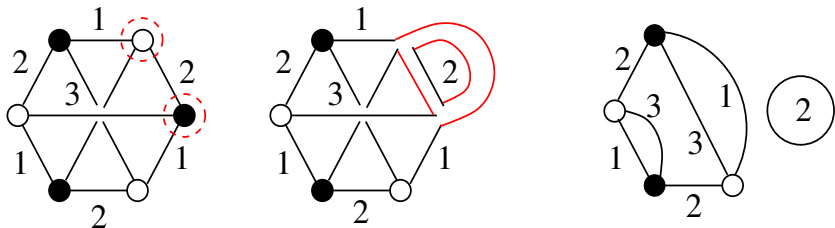


$$\sum_{a_1, \dots, a_d} \int [dTd\bar{T}] \frac{\partial}{\partial \bar{T}_{a_1 \dots a_d}} \left((B \setminus v)_{a_1 \dots a_d} e^{-N^{d-1} S(T, \bar{T})} \right) = 0$$

Operations on bubbles

Contraction

Cut 2 vertices v, \bar{v} and connect the lines pairwise

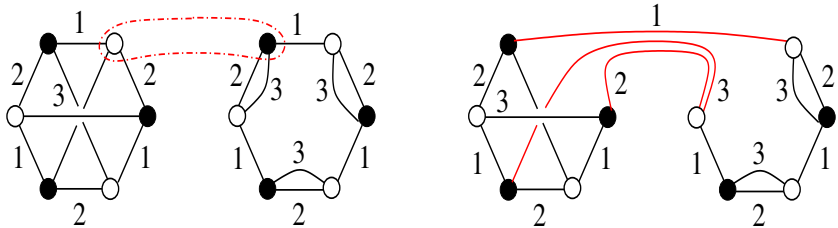


$$\sum_{a_1, \dots, a_d} \frac{\partial (B \setminus v)_{a_1 \dots a_d}}{\partial \bar{T}_{a_1 \dots a_d}} = \sum_{\bar{v} \in B} B \setminus (v, \bar{v})$$

Operations on bubbles

Gluing

Take 2 open bubbles and glue them



$$\sum_{a_1, \dots, a_d} (B_1 \setminus v_1)_{a_1 \dots a_d} (B_2 \setminus \bar{v}_2)_{a_1 \dots a_d} = (B_1 \setminus v_1) \cdot (B_2 \setminus \bar{v}_2)$$

Schwinger-Dyson Equations

For any open bubble $B \setminus v$, at all orders

$$\frac{1}{N^{d-1}} \sum_{\bar{v} \in B} \underbrace{\langle B \setminus (v, \bar{v}) \rangle}_{\text{contraction of } B} - \langle B \rangle - \sum_{i \in I} t_i \sum_{\bar{v} \in B_i} \underbrace{\langle (B \setminus v) \cdot (B_i \setminus \bar{v}_i) \rangle}_{\text{gluing with } B_i} = 0$$

- ▶ $\langle B \rangle = \partial_{t_B} \ln Z$
Translate SDEs as constraints $L_{(B \setminus v)} Z = 0$ with differential operators $L_{(B \setminus v)}$
- ▶ Infinite-dimensional Lie algebra

$$[L_{B_1 \setminus v_1}, L_{B_2 \setminus v_2}] = \sum_{\bar{v}_1 \in B_1} L_{((B_1 \setminus \bar{v}_1) \cdot (B_2 \setminus v_2)) \setminus v_1} - \sum_{\bar{v}_2 \in B_2} L_{((B_1 \setminus v_1) \cdot (B_2 \setminus \bar{v}_2)) \setminus v_2}$$

- ▶ $d = 2$, all bubbles are loops labeled by number of vertices p ,

$$[L_{p_1}, L_{p_2}] = (p_1 - p_2) L_{p_1 + p_2}$$

Gaussian models

- ▶ For a real variable ϕ , covariance λ ,

$$\langle \phi^{2p} \rangle = \sum_{\text{Wick pairings}} \lambda^p$$

- ▶ For matrix M at large N , not all pairing have the same scaling!

$$\langle \text{tr } M^{2p} \rangle = \sum_{\text{Wick pairings}} N^{2-2h} \lambda^p = \sum_{\text{planar pairings}} N \lambda^p = N \lambda^p c_p$$

where c_p Catalan number.

- ▶ For tensor at large N , bubble with $2p$ vertices, only one melonic contraction

$$\langle B_{2p} \rangle = \sum_{\text{Wick pairings}} N^{d-\omega} \lambda^p = \sum_{\text{melonic pairings}} N \lambda^p = N \lambda^p$$

Solution and Universality at large N

Revisiting random tensor models at large N via the Schwinger-Dyson equations,
V. Bonzom, JHEP (2012)

Schwinger-Dyson equations have a single perturbative solution

- ▶ For any two melonic bubbles with $2p$ vertices, $\langle B \rangle = \langle B' \rangle \equiv N G_p$
- ▶ Physical solution

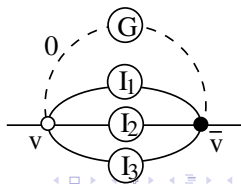
$$G_p = G^p, \quad \text{for} \quad 1 - G - \sum_{i \in I} p_i t_i G^{p_i} = 0$$

Algebraic equation, and take branch $G = 1$ when $t_i = 0$

- ▶ $p = 1$ gives the full 2-point function $\langle T \cdot \bar{T} \rangle = N G$

Models are **Gaussian** at large N

- ▶ l_1, l_2, l_3 : melonic insertions
- ▶ G : insertion of full 2-point function

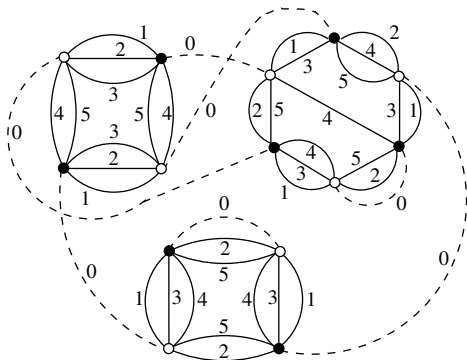


Is it possible to retain more than melonic graphs?

New $1/N$ expansions, V. Bonzom, JHEP (2012)

Probe sub-graphs

- ▶ Split the set of colors into two sets $\{1, \dots, k\}$, $\{k+1, \dots, d\}$
- ▶ Compare degree of Feynman graphs with degrees of subgraphs with colors $\{0, 1, \dots, k\}$ and $\{0, k+1, \dots, d\}$

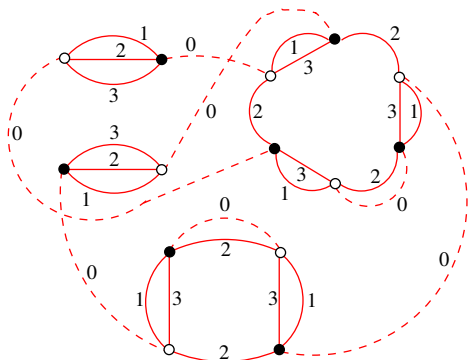


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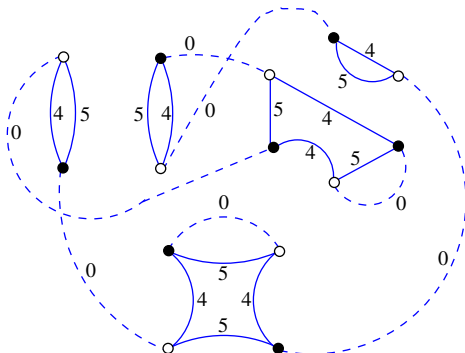


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$$\omega(G) = \omega_{01\dots k} + \omega_{0k+1\dots d} + (l - b)$$

- ▶ Amplitude of a graph

$$A(G) \sim N^{-\omega(G)} \lambda^{l-b} = N^{-\omega_{01\dots k} - \omega_{0k+1\dots d}} \left(\frac{\lambda}{N}\right)^{l-b}$$

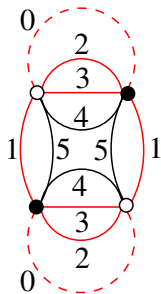
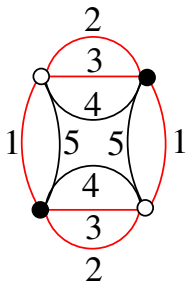
- ▶ Suggests to scale $\lambda = N \kappa$

$$\int [dT d\bar{T}] \exp -\frac{N^{d-2}}{\kappa} S(T, \bar{T}) \quad \text{i.e. } \alpha = d - 2$$

- ▶ Danger: melonic graphs become unbounded
Rescaling melonic bubbles is sufficient
- ▶ It re-shuffles the $1/N$ expansion
Add new large N contributions

Large N : Allows more than melonic bubbles!

- ▶ sub-bubbles $\{1 \dots, k\}$ melonic
- ▶ sub-bubbles $\{k + 1 \dots, d\}$ melonic
- ▶ subgraphs $\{0, 1 \dots, k\}$ melonic
- ▶ subgraphs $\{0, k + 1, \dots, d\}$ melonic



- ▶ Still Gaussian model

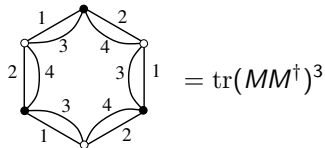
What else?

Matrix models are non-Gaussian!

- ▶ Introduce *fat indices*

$$A = (a_1, a_3), \quad B = (a_2, a_4)$$

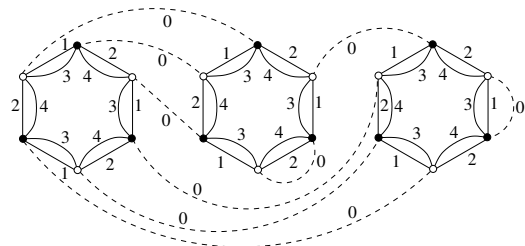
$$T_{a_1 a_2 a_3 a_4} = M_{AB}$$



$$= \text{tr}(MM^\dagger)^3$$

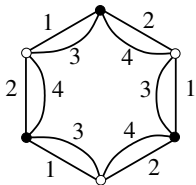
- ▶ A model that generates triangulations of the 4-sphere at large N

$$\int [dT d\bar{T}] \exp -N^2 \left(T \cdot \bar{T} + \text{Diagram} \right)$$

$$= \text{large } N \sum_{\text{Triangulations 4-sphere}}$$


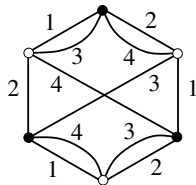
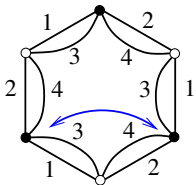
Extend the properties of that model

- ▶ Consider bubbles with a single loop (1,2) and a single loop (3,4)



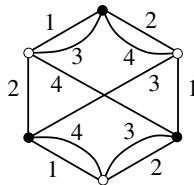
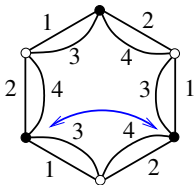
Extend the properties of that model

- ▶ Consider bubbles with a single loop (1,2) and a single loop (3,4)
- ▶ Draw the loop (3,4) after a permutation on the black vertices

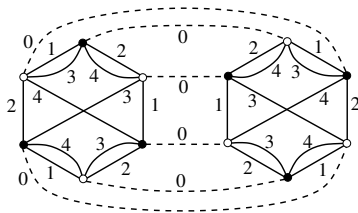
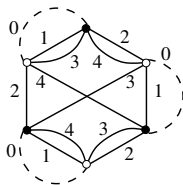


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Non-Gaussian contributions



Conclusion

Tensor models provides control on random triangulations

- ▶ Generate random colored triangulations from bubbles
- ▶ Large N : maximize number of faces
- ▶ $d = 2$, planar graphs ; $d \geq 3$, melonic or have melonic subgraphs
- ▶ In probability language: Gaussianity
- ▶ However non-Gaussian models do exist, in dimension 4!

Further explore

- ▶ Today talk: ways to probe post-melonic world
- ▶ Another way: double scaling
- ▶ Rich combinatorics
- ▶ There is a closed system of equations: Schwinger-Dyson, but gets more complicated
- ▶ Preliminary investigation: study families of bubbles