

Tensor Models/Group Field Theories: An overview

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Outline

- 1 Introduction: Overview of Tensor Models/Group Field Theory
- 2 Tensor Models & Large N expansion
 - Invariant Tensor Models
 - Colored Graphs, Invariants, Jackets and all that...
 - Large N limit and critical behavior
 - More results
- 3 Tensorial Group Field Theory: Renormalization Program
 - Defining Tensor-like QFTs
 - Renormalization
 - Overview of renormalizable models
- 4 TM Combinatorics
 - Gaussian Moments and Meanders
 - Counting TM Observables/Permutation-TFT
- 5 Conclusion: Open Questions/Next Challenges

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Generalities: Quantizing Gravity

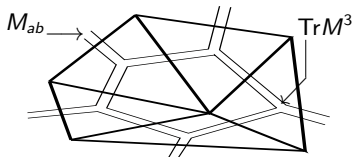
- Several issues. The Einstein theory of gravity is perturbatively divergent and nonrenormalizable [DeWitt PR '67, Goroff & Sagnotti, NPB '86].
- Remarkable Sakharov's idea: "An induced theory of gravity" [Sakharov, '67, SPD 68, Visser MPLA '02].
- Alternative scenarii (coupling gravity to other fields, Asymptotic Safety, ...) and more "daring" scenarii (extra-dimensions, susy, background field independent methods, ...).
- Focus on *Discrete* methods. Mid 80's: In particular, *Matrix Models* [Di Francesco et al., PR '95] prove to be a solvable framework and concrete realization of an "emergent gravity" scenario.
- Matrix Models and Random 2D geometry: To "replace" the sum over topologies and geometry of a 2D surface by a sum over random triangulations of surfaces.

From Matrices ...

Matrix Models (MM): Probability measures for matrices M of *large size* N and describe 2D gravity. Archetype:

$$Z_{\text{matrix}} = \int dM e^{-\frac{1}{2} \text{Tr} M^2 + \frac{\lambda}{\sqrt{N}} \text{Tr} M^3} = e^{Z_{2DQG}}$$

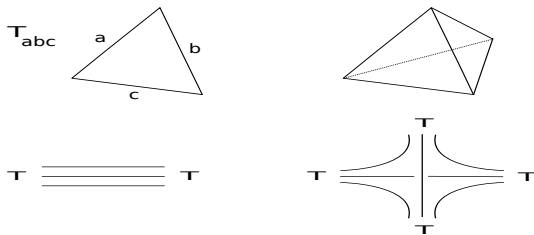
- A triangulated surface \equiv A Feynman ribbon graph:



- 't Hooft's Large N limit: Planar graphs' sector \equiv surfaces of genus 0 (can be counted like trees).
- Stat. Mech.: \exists phase transition ($N \rightarrow \infty$; $\lambda \rightarrow \lambda_c$) \rightsquigarrow a continuum limit (infinitely refined Riemann surfaces) as a 2D theory of gravity (Liouville + CF matter).

... to Tensors

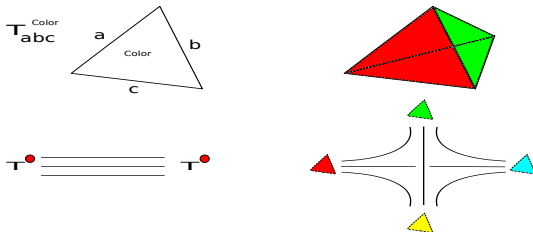
- Tensor Models (TM) of rank D : Tool for randomizing geometry in dimension D
- Basic building blocks $(D - 1)$ -simplexes & Interaction forms a D -simplex;
- For e.g. in 3D:



- Some results [Ambjorn et al. '91, Sasakura '91, Gross '92, Boulatov-Ooguri '92]
 - Ambjorn et al., '91', numeric phase transition;
 - Boulatov-Ooguri-type TM for 3D-4D simplicial gravity: (Topological/Lattice Field Theory like theory) '92/'93'; Related to Ponzano-Regge partition function 3D;
 - Loop Quantum Gravity Connection ([Reisenberger-Rovelli, '00, Freidel, '05, Oriti, '06]) and "birth" of Group Field Theory.
- Lack of $1/N$ expansion \Rightarrow all nice exact results of MM cannot be extended to TM.
- TM need improvement(s): $C + DT$'s, Boulatov-Ooguri models and GFT's;

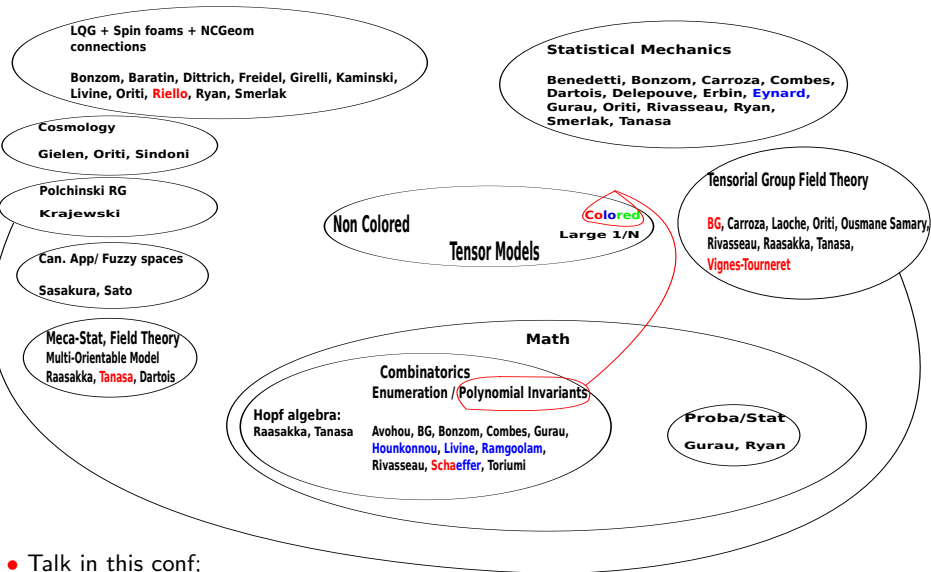
Colored Tensor Models

- '10 Gurau's $1/N$ expansion for colored TM [Gurau, AHP, '11]
- 3D:



- triangulate better objects (pseudo-manifolds) [Gurau, CMP '11]
- Leading graphs triangulate only spheres in any D [Gurau, AHP '11]
- have computable critical exponent [Bonzom, Gurau, Riello, Rivasseau, NPB, '11];
- with possible matter fields [Bonzom et al., PRD '12 ; Benedetti et al, NPB '12];
- could lead to extension of the Virasoro-DeWitt algebra [Gurau, NPB '11];
- underlie universal theory for general (ie unsymmetrized) tensors [Gurau, '11];
- Existence of a double-scaling limit [Dartois et al. '13, Gurau & Schaeffer, '13]
- Expand a natural $U(N)^{\otimes D}$ invariance [Gurau, '12, Bonzom et al. '12]
- Define renormalizable field theories;

TM People



- Talk in this conf;
- Recent Honorable Guest Contributors;
- Ordinary Contributors (running the program with usual/stressful deadlines);

Foundation questions:

- Tensor: Only a discretization tool or a “real” quanta of some “thing” ?

Big bang is identified with **geometrogenesis**, ie emergence of classical space-time through one or several **phase transitions**. Pre-space (analog of space-time before condensation) is treated as a **physical transplanckian early phase of the universe** (not just as a mathematical trick) [Orti, '06, Konopka '08].

- What if $N = \infty$ from the beginning ? Several MM will be already divergent, worse are the behavior of TM.

- Divergences \Rightarrow Need of **Renormalization**. Basic axioms for a QFT with Tensors ? [Rivasseau: The Tensor Track '11, '12, '13] **Renormalization group is a guiding/selecting thread** in models' space.

- Does it work for matrices ? The Grosse-Wulkenhaar model: Renormalizable and Asymptotically Safe model (without any extra-symmetry required) and induces important NEW developments in Field Theory. Closed equations might be extended for higher rank tensor [Ousmane Samary, '13] but the resolution is truly challenging.

- Motto: “Randomizing Geometry = Quantum Gravity”
- Goal: Achieve a universal scenario for an emergent spacetime through one or several phase transitions.
- Tools/Methods/Motivations: At the crossroad of Matrix models (large N limit), Group Field Theory, QFT/Constructive renormalization, Proba/Stats.

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Tensor Models [Gurau, '10, '11]

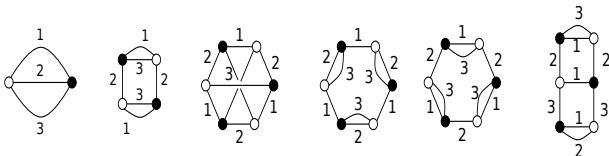
- Study of probability measures of random tensor spaces + Geometric/Topological/Physical inputs.
- A covariant complex tensor T_{p_1, \dots, p_d} with transformation rule

$$T_{p_1, \dots, p_d}^U = \sum_{q_k} U_{p_1 q_1}^{(1)} \dots U_{p_d q_d}^{(d)} T_{q_1, \dots, q_d}, \quad U^{(a)} \in U(N_a) \quad (1)$$

- G/T/Physics input: T is viewed as a $(d-1)$ -simplex.
- **Tensor Invariance** for defining the interactions

$$S_b^{\text{int}}(T, \bar{T}) = \text{Tr}_b(T \dots T) = S_b^{\text{int}}(T^U, \bar{T}^U) \quad (2)$$

b a colored graph encoding the contraction pattern; S_b^{int} "is" a gluing of simplexes and represents a d -simplex. Ex:



(courtesy of Gurau)

Action/Partition function

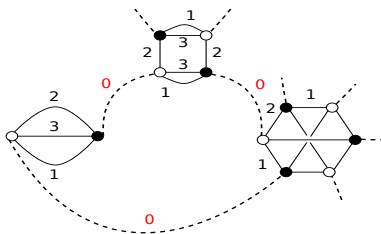
- Invariant action:

$$S[T, \bar{T}, \{\lambda_{\mathbf{b}}\}_{\mathbf{b}}] = \sum_{\mathbf{b}} \lambda_{\mathbf{b}} N^{-\omega(\mathbf{b}, d)} S_{\mathbf{b}}^{\text{int}}(T, \bar{T}), \quad \omega(\mathbf{b}, d) \geq 0 \quad (3)$$

- Partition function

$$Z[\{\lambda_{\mathbf{b}}\}_{\mathbf{b}}] = \int \prod_{p_i} [dT_{p_1, \dots, p_d} d\bar{T}_{p_1, \dots, p_d}] e^{-N^{D-1} S[T, \bar{T}, \{\lambda_{\mathbf{b}}\}_{\mathbf{b}}]} \quad (4)$$

- Expanding the partition function: Feynman graphs

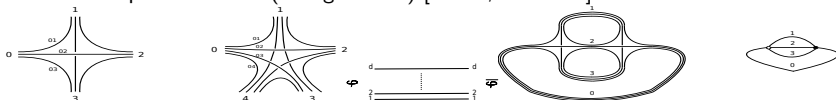


- But what is the analogue of matrix models' ribbon graphs ?

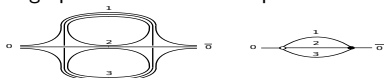
Graph Anatomy:
On the case of a Colored Tensor Graph

Colored Tensor Graphs I

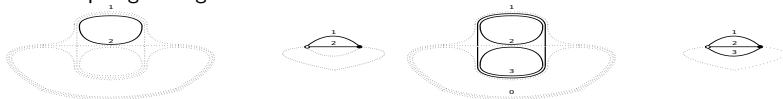
- Colored bi-partite models (triangulations) [Gurau, CMP '11]:



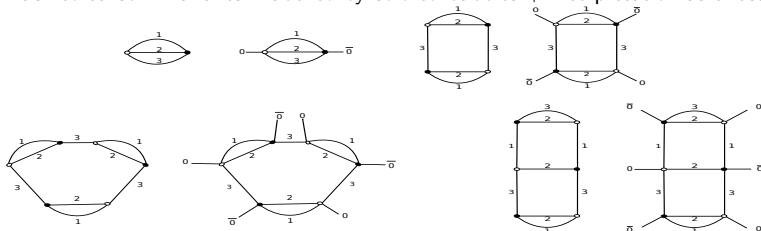
A graph can be closed or open



Other topological ingredients: Faces and bubbles

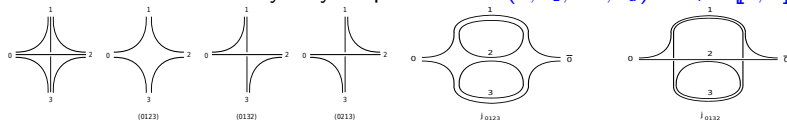


Back to tensor invariants: Labelled by colored bubbles + Interpretation as effective vertices.



Colored Tensor Graphs II

Jacket: Ribbon indexed by a cyclic permutation $(0, a_1, \dots, a_d)$ of $a_i \in \llbracket 1, d \rrbracket$



Boundary graph $\partial\mathcal{G}$: Rank $d - 1$ graph encoding the boundary of the simplicial complex associated with \mathcal{G} .

Large N limit and critical behavior

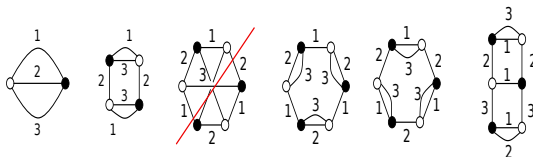
- Precise action: “Single” trace-like with weights allows to obtain, for a connected graph \mathcal{G} ,

$$A(\mathcal{G}) = N^{d - \frac{2}{(d-1)!} \omega(\mathcal{G})}, \quad \omega(\mathcal{G}) = \sum_{J \subset \mathcal{G}} g_J \quad (5)$$

- $\omega(\mathcal{G})$ called the degree of \mathcal{G} . [Gurau '10, Bonzom, Gurau, Riello, Rivasseau, '12]

Model Type	Matrix	Rank d Colored Tensor
Expansion in	$g(\mathcal{G})$ genus	$\omega(\mathcal{G})$ degree
Amplitudes	$A(\mathcal{G}) = N^{2-2g(\mathcal{G})}$	$A(\mathcal{G}) = N^{d - \frac{2}{(d-1)!} \omega(\mathcal{G})}$
Leading graphs	Planar $\equiv S^2$	Melon $\equiv S^d$
$F_{\text{Lead}} \sim (\lambda_c - \lambda)^{2-\gamma}$	$\gamma_{\text{planar}} = -\frac{1}{2}$	$\gamma_{\text{melonic}} = \frac{1}{2}$

- “Melons” ?



- Phase transition \rightsquigarrow Branched Polymer phase [Gurau & Ryan, '13]: The new phase of Colored TM defines a Continuous Random Tree in the sense of Aldous (same (Hausdorff, Spectral dimension) = $(2, 4/3)$).

More results

- The Ising model in dimension D [Bonzom et. al, '12] on random tensor graphs (lattices): No phase transition at finite coupling (agrees with numerics); To force a phase transition modify the models;
- Universality [Gurau, '12]: Measures on random tensors converge (in distribution, in a precise sense) to the Gaussian measure at the large N limit.
- Multi-critical behavior on random (spherical) lattices and (hard) dimers [Bonzom, '12; Bonzom & Erbin, '12]: Coupling tensors interpreted as dimer activities; Phase transition, $\gamma_m = 1 - \frac{1}{m}, m \geq 2$, \leadsto hard dimers on branched polymers.
- Double scaling in a T^4 model [Dartois et al. '13; Gurau & Schaeffer, '13]:

New double scaling limit for TM $N \rightarrow \infty, \quad \lambda \rightarrow \lambda_c, \quad N^\alpha(\lambda - \lambda_c) = \text{const.} \leadsto \text{Polymers} \quad (6)$

Using constructive techniques (loop vertex expansion), the series is Borel summable if $d < 6$.

- Beyond perturbations/Constructive techniques [Gurau, '13; Delpouve et al., '14]: The meaning of the $1/N$ -expansion, reaching the critical point by analytic continuation $\lambda \rightarrow \lambda_c < 0$; Borel summability of T^4 quartic models.

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Defining Tensor-like QFTs

“... one can define TGFT's with tensorial interactions and a soft breaking of the tensorial invariance of their propagator. ... + desirable features + ... + ...” (V. Rivasseau, The Tensor Track III, 1311.1461)

- **Simple** TM: A complex tensor T_{p_1, \dots, p_d} with “Tensor Invariance” for defining the interactions $S_b^{\text{int}}(T, \bar{T}) = \text{Tr}_b(T \dots T) = S_b^{\text{int}}(T^U, \bar{T}^U)$ and Kinetic term

$$S^{\text{kin}}(T, \bar{T}) = \sum_{p_s} \bar{T}_{p_1, \dots, p_d} \left(\sum_s (p_s)^{2a} + \mu \right) T_{p_1, \dots, p_d} \quad (7)$$

$0 < a \leq 1$; $(p_s^2)^a \equiv (\Delta)^a$ and $a \leq 1$ might be useful to recover O.S. positivity axiom (Rivasseau, Tensor Track III, '13).

- **Gauge invariant models**: Imposing constraints on T (Carrozza, Oriti, Ousmane Samary, Rivasseau, Vignes-Tourneret).
- Summing over arbitrary high momenta may imply divergent amplitudes.



Renormalization à la V. Rivasseau

Base Cut off $\Lambda \rightarrow \infty$

$(g_n)_{n \geq 0}$ $\Delta \rightarrow \infty$ $\xrightarrow{\text{sans finie}} 1949$

$f = Ag$ $g(\Lambda) = g_n$ $f = g_n$ $\Lambda \rightarrow \infty$

hate une base \rightarrow $\int_{\Lambda}^{\Lambda} d\mu$

Wilson 1973* \downarrow \rightarrow $\text{cutoff rules (renormalization)}$

$C_i = \int_{\Lambda}^{\Lambda} \frac{e^{-\alpha k}}{1 - z_i} - \text{eff} = \text{ren.}$

$S = \sum_{n=1}^{\infty} \frac{\Lambda^{n-1}}{n!} A_i$

$\sum_{n \in \mathbb{N}} a_n \Lambda^n = \sum_{n_i} a_{n_i} \Lambda^{n_i}$

$\sum a_{n_i} \Lambda^{n_i} y_{n_i}$ $y = A \log x$

$\Lambda \rightarrow \infty$

$|I_G(q) - \text{CT}_{G(q)}^{\text{P} \geq q}| \xrightarrow{2}$

théorie effective

$\text{localité} \Leftrightarrow \text{décomposition en edulle telle que}$

BPHZ

$\text{Reuermann Bogdanov Parashuk Hepp} \rightarrow \text{amplitude finie}$

$\text{Zimmerman} \rightarrow \text{solution de la renormalisation sous-traction par forêt}$

$I_G = \sum_{F \subset G} \frac{(-1)^{|F|}}{|F|!} I_{G-F}^{\text{Base}}$

Renormalization of TGFTs

in other words: QFT Renormalization is the “intricate” combination of a Multi-Scale Analysis, a Power-Counting Theorem and a Locality Principle.

- Consider $T : U(1)^d \rightarrow \mathbb{C}$ and its Fourier components T_{p_1, \dots, p_d}

$$Z = \int d\nu_C(T, \bar{T}) e^{-S^{\text{int}}(T, \bar{T})}, \quad S^{\text{int}}(T, \bar{T}) = \sum_{\mathbf{b}} \lambda_{\mathbf{b}} \text{Tr}_{\mathbf{b}}(T \dots T) \sim \sum_n \lambda_n \phi^{2n} \quad (8)$$

- Multi-Scale Analysis: Slice decomposition: $C = \sum_{i=0}^{\infty} C_i$

$$\begin{aligned} C_i &= \int_{M^{-2i}}^{M^{-2(i-1)}} d\alpha e^{-\alpha(\sum_{s=1}^d p_s^{2a} + \mu)}, \quad \forall i \geq 0, M > 1 \\ &\leq KM^{-2i} e^{-\delta M^{-i}(\sum_{s=1}^d p_s^{2a} + \mu^2)} \quad K > 0, \delta > 0 \end{aligned} \quad (9)$$

High i probes high p ; Cut-off $C^\Lambda = \sum_{i=0}^\Lambda C_i$;

- Power-Counting Theorem: $|A_{\mathcal{G}, \mu}| \leq K^n M^{\omega_d(\mathcal{G})}$ where the divergence degree of a graph \mathcal{G} is given by

$$\omega_d(\mathcal{G}) = -A_d(\omega(\mathcal{G}_{\text{color}}) - \omega(\partial\mathcal{G})) - (C_{\partial\mathcal{G}} - 1) - B_d[N_{\text{ext}} - C_d] - \frac{1}{2}D_d \cdot V. \quad (10)$$

- Locality Principle: Melons with melonic boundary are dominant $\omega(\mathcal{G}_{\text{color}}) = 0$; $\omega(\partial\mathcal{G}) = 0$

Results: Several Renormalizable Models

TGFT (type)	G_D	$\phi^{k_{\max}}$	d	a	Renormalizability	UV behavior
	$U(1)$	ϕ^6	4	1	Just-	AF
	$U(1)$	ϕ^3	3	$\frac{1}{2}$	Just-	AF
	$U(1)$	ϕ^6	3	$\frac{1}{2}$	Just-	AF
	$U(1)$	ϕ^4	4	$\frac{3}{4}$	Just-	AF
	$U(1)$	ϕ^4	5	1	Just-	AF
	$U(1)^2$	ϕ^4	4	1	Just-	AF
	$U(1)$	ϕ^{2k}	3	1	Super-	-
gi-	$U(1)$	ϕ^4	6	1	Just-	AF
gi-	$U(1)$	ϕ^6	5	1	Just-	AF
gi-	$SU(2)^3$	ϕ^6	3	1	Just-	AF
gi-	$U(1)$	ϕ^{2k}	4	1	Super-	-
gi-	$U(1)$	ϕ^4	5	1	Super-	-
Matrix	$U(1)$	ϕ^{2k}	2	$\frac{1}{2}(1 - \frac{1}{k})$	Just-	$(k = 2, \text{AS}^{(\infty)}); (k = 3, \text{LG})$
Matrix	$U(1)^2$	ϕ^{2k}	2	$1 - \frac{1}{k}$	Just-	$(k = 2, \text{AS}^{(1)}); (k = 3, \text{LG})$
Matrix	$U(1)^3$ or $SU(2)$	ϕ^6	2	1	Just-	LG
Matrix	$U(1)^3$ or $SU(2)$	ϕ^4	2	$\frac{3}{4}$	Just-	$\text{AS}^{(1)}$
Matrix	$U(1)^4$	ϕ^4	2	1	Just-	$\text{AS}^{(1)}$
Matrix	$U(1)$	ϕ^{2k}	2	$\frac{1}{2}$	Super-	-
Matrix	$U(1)^2$	ϕ^{2k}	2	1	Super-	-

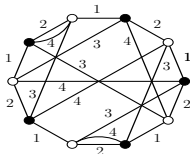
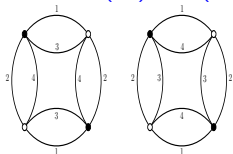
Table: Updated list of tensorial renormalizable models and their features (AF \equiv asymptotically free; LG \equiv existence of a Landau ghost; $\text{AS}^{(\ell)}$ asymptotically safe at ℓ -loops).

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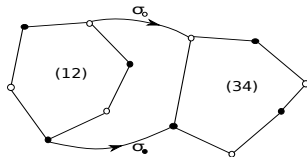
Meanders [Bonzom & Combes, '13]

- Question: Computation of polynomials $P_b(T, \bar{T})$ expectation values for the Gaussian measure, in the large N but including as much corrections;
- \exists nice answer: In $d = 4$, if \mathbf{b} has two unique faces labelled by a couple of colors, for instance, (12) and (34).



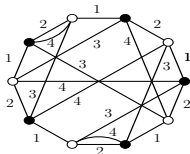
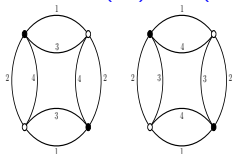
courtesy of Bonzom

- $\mathbf{b}_{\sigma_o, \sigma_\bullet}$ labelled by 2 permutations such that



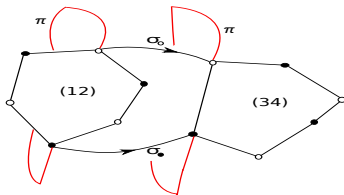
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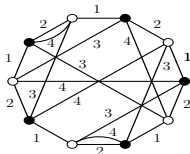
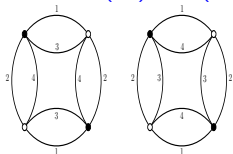
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- b_{σ_o, σ_e} labelled by 2 permutations such that



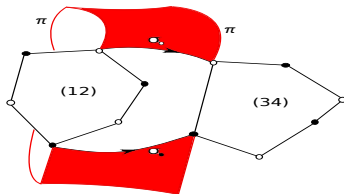
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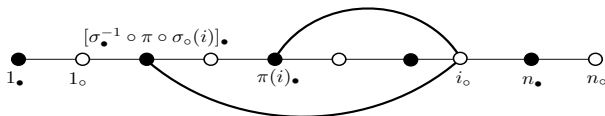
courtesy of Bonzom

- $b_{\sigma_o, \sigma_\bullet}$ labelled by 2 permutations such that



Meanders

- Using a nice labelling of vertices of both cycles \rightsquigarrow Meanders



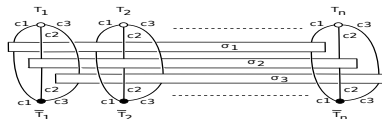
- $\mathbf{b}_{\sigma_o, \sigma_\bullet}$ such that

$$\begin{aligned}
 \langle P_{\sigma_o, \sigma_\bullet} \rangle &= \frac{1}{Z} \int e^{-N^2 T \cdot \bar{T}} P_{\sigma_o, \sigma_\bullet} dT d\bar{T}, = \sum_{\text{paring } \pi} N^{\Omega(\sigma_o, \sigma_\bullet, \pi)} \\
 &= N^2 \sum_{\pi} N^{-2g_{012; \pi} - 2g_{034; \pi}} =_{\text{large } N} N^2 |\mathcal{M}_{\sigma_o, \sigma_\bullet}|
 \end{aligned} \tag{11}$$

- Subleading contributions can be evaluated in terms of irreducible meanders systems associated with polynomials labelled by stabilized-interval-free permutations.

Counting TM Observables/Permutation-TFT [JBG & Rangoolam, '13]

- Rank $d = 3$: Determination of possible graphs $\text{Tr}_b(T, \bar{T})$ amounts to count 3-uples



$$(S_n \times S_n \times S_n) \ni (\sigma_1, \sigma_2, \sigma_3) \sim (\gamma_1 \sigma_1 \gamma_2, \gamma_1 \sigma_2 \gamma_2, \gamma_1 \sigma_3 \gamma_2) \quad (12)$$

- Counting points in the double coset

$$\mathcal{S}_3(n) = \text{Diag}(S_n) \backslash (S_n \times S_n \times S_n) / \text{Diag}(S_n). \quad (13)$$

- Burnside's lemma: Permutation-TFT formulation !

$$Z_3(n) = \frac{1}{(n!)^2} \sum_{\gamma_{1,2} \in S_n} \sum_{\sigma_i \in S_n} \left(\delta(\gamma_1 \sigma_1 \gamma_2 \sigma_1^{-1}) \delta(\gamma_1 \sigma_2 \gamma_2 \sigma_2^{-1}) \delta(\gamma_1 \sigma_3 \gamma_2 \sigma_3^{-1}) \right) \quad (14)$$

$$a \rightarrow \sigma_1, \quad b \rightarrow \sigma_2, \quad c \rightarrow \sigma_3,$$



- Solve one delta, say σ_1 (fixing the gauge), $\sigma_1^{-1} \sigma_i = \tau_{i-1}$, inserting $\tau_0 / \tau_0 \tau_1 \tau_2 = 1$ and get

$$Z_3(n) = \frac{1}{n!} \sum_{\gamma, \tau_0 \in S_n} \sum_{\tau_{1,2} \in S_n} \left(\delta(\gamma \tau_1 \gamma^{-1} \tau_1^{-1}) \delta(\gamma \tau_2 \gamma^{-1} \tau_2^{-1}) \delta(\tau_0 \tau_1 \tau_2) \right) \quad (15)$$

counts equivalent classes of branched covers of $S^2 \setminus \{3\bullet\}$! (I dunno yet why but this is important 😊)

Outline

- 1 Introduction: Overview of Tensor Models/Group Field Theory
- 2 Tensor Models & Large N expansion
 - Invariant Tensor Models
 - Colored Graphs, Invariants, Jackets and all that...
 - Large N limit and critical behavior
 - More results
- 3 Tensorial Group Field Theory: Renormalization Program
 - Defining Tensor-like QFTs
 - Renormalization
 - Overview of renormalizable models
- 4 TM Combinatorics
 - Gaussian Moments and Meanders
 - Counting TM Observables/Permutation-TFT
- 5 Conclusion: Open Questions/Next Challenges

Open questions/Next Challenges

- Tensor models generalize matrix models and have, in a sense (after restrictions) better behavior: The double scaling limit is possible for T^4 models and the series is Borel summable whereas the BP phase persists.
 - ~ Open question: Determine another summable family of subleading graphs which may change this phase (triple/multiple scaling). Rinse, Repeat....
 - ~ Challenge: The new phase should be not model dependent (universality).
- Field Theory: There exist several models which can be renormalizable and AF.
 - ~ Open questions: • Determine conditions to truncate this space of model (Renormalizability and AF seems to be too weak). • Prove the OS positivity axiom and first steps towards a 4D Minkowskian spacetime; • There exist certainly different types of TGFT than the one present here and Renormalization analysis works as well: Start by a matrix-vector coupling with model covariance:

$$S[M, \phi] = \text{Tr} \left(M_{ab}(a + b + \mu)M_{ab} + \phi^a(\mu_\phi + a)\phi^a + M_{ab}(a + b + \mu)\phi^a\phi^b \right) + V(M, \phi) \quad (16)$$

- ~ Challenge: • Hunting nonperturbative and universal effects (FRG methods, recovering phase diagrams) • Make the rank d of tensors or the group dimension $\dim G$ dynamical.
- Combinatorics: • Elucidate the correspondence between equivalent classes of S^2 branched covers and the counting TM observables; • Generic formula for the TM correlation functions ? • How the enumeration of TI are useful for other conjectures in Math (e.g. Gromov's on sphere triangulations) ? • Polynomials Invariant generalizing Tutte/Bollobas Riordan: Universality theorem for a 3D polynomial invariant. New polynomial invariants issued from Tensor Invariants FT (wip with R. Toriumi).