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(Non)-geometry and non-geometric fluxes

David ANDRIOT

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(Partial) reviews: [arXiv:0708.3984](#), [1106.4015](#), [1303.0251](#)

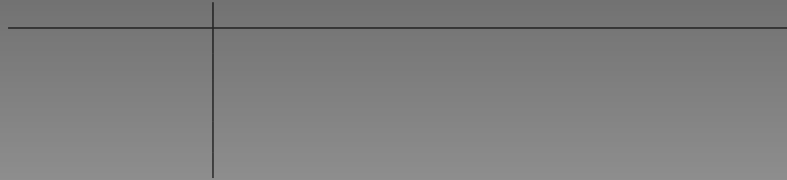
[arXiv:1306.4381](#), [1402.5972](#) by D. A. and André Betz

Quantum Gravity in Paris,
LPT Orsay and IHES, 18/03/2014, France

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A standard scenario to relate to string theory to 4D physics



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	theory
world-sheet (2D)	Superstring (CFT)

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A standard scenario to relate to string theory to 4D physics

	theory	NSNS sector
world-sheet (2D)	Superstring (CFT) $\mathcal{L}_\Sigma \sim (g_{mn} + b_{mn})(\mathcal{X}) \partial_- \mathcal{X}^m \partial_+ \mathcal{X}^n$	

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10D target space	Supergravity	

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at all levels, but respect a priori none of the steps

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at all levels, but respect a priori none of the steps
 \hookrightarrow review them from various perspectives,
give rise to interesting (non-standard) geometric structures
- 10D β -supergravity: restores standard geometric description
 \hookrightarrow allows compactification \Rightarrow (new?) phenomenology

Different perspectives on non-geometry

Four-dimensional non-geometric fluxes

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Four-dimensional non-geometric fluxes

NSNS sector of four-dimensional gauged supergravity

Super/scalar potential depends on fluxes (constants),

e.g. H_{abc} , $f^a{}_{bc}$, but also non-geometric fluxes: $Q_a{}^{bc}$, R^{abc}

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Components of embedding tensor

\Leftrightarrow some “structure constants” in gauging algebra

hep-th/0508133 by J. Shelton, W. Taylor, B. Wecht

hep-th/0210209, hep-th/0512005 by A. Dabholkar, C. Hull

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\Rightarrow specific terms in 4D potential (\checkmark for phenomenology!)

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Terms/fluxes related by 4D T-duality chain

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10D origin: H , f from compactification; what about Q and R ?

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\leftrightarrow answers from β -supergravity

Ten-dimensional non-geometry

Original idea of non-geometry

[hep-th/0208174](#) by S. Hellerman, J. McGreevy, B. Williams

[hep-th/0210209](#) by A. Dabholkar, C. Hull

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A (target) space, divided in patches

Fields glue on overlaps:

diffeomorphisms, gauge transformation

(point-like symmetries)

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\leftrightarrow use stringy symmetries for gluing

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Away from standard geometry

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Fields look ill-defined:

not single-valued, global issues

Simple example: one T-duality: circle of radius $R \rightarrow \frac{1}{R}$

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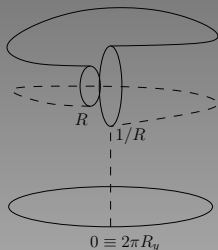
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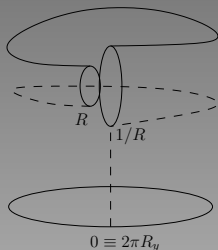
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No compactification on such spaces...



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Famous toroidal example: NSNS sector:

$$T^3 + H_{123} \xleftrightarrow{T_1} \text{twisted torus } (f^1_{23}) \xleftrightarrow{T_2} \text{non-geometric config.}$$

[hep-th/0211182](#) by S. Kachru, M. B. Schulz, P. K. Tripathy, S. P. Trivedi

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$$g = f_0 \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & \frac{1}{f_0} \end{pmatrix}, \quad b = f_0 \begin{pmatrix} 0 & -Kz^3 & 0 \\ Kz^3 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}, \quad f_0(z^3) = \frac{1}{1 + (Kz^3)^2}$$

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2 patches for $S^1_{z^3}$, proof that no diffeo. to glue g on overlaps
+ fields glue with T-duality element $\in O(2,2) \Rightarrow$ non-geometry

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Other example: NS-branes (NSNS backgrounds)

$NS5\text{-brane} \xleftrightarrow[+T\text{-d.}]{\text{smearing}} KK\text{-monopole} \xleftrightarrow[+T\text{-d.}]{\text{smearing}} 5^2_2 \text{ or } Q\text{-brane}$

[arXiv:1004.2521](#), [1209.6056](#) by J. de Boer and M. Shigemori

[arXiv:1303.1413](#) by F. Hassler and D. Lüst, [arXiv:1402.5972](#) by D. A. and André Betz

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Comparison to the 4D T-duality chain \Rightarrow Q -flux ?!

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$$g = f_0 \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & \frac{1}{f_0} \end{pmatrix}, \quad b = f_0 \begin{pmatrix} 0 & -Kz^3 & 0 \\ Kz^3 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}, \quad f_0(z^3) = \frac{1}{1 + (Kz^3)^2}$$

2 patches for $S^1_{z^3}$, proof that no diffeo. to glue g on overlaps
+ fields glue with T-duality element $\in O(2,2) \Rightarrow$ non-geometry

[arXiv:1402.5972](#) by D. A. and André Betz

Other example: NS-branes (NSNS backgrounds)

$NS5\text{-brane} \xleftrightarrow[+T\text{-d.}]{\text{smearing}} KK\text{-monopole} \xleftrightarrow[+T\text{-d.}]{\text{smearing}} 5^2_2 \text{ or } Q\text{-brane}$

[arXiv:1004.2521](#), [1209.6056](#) by J. de Boer and M. Shigemori

[arXiv:1303.1413](#) by F. Hassler and D. Lüst, [arXiv:1402.5972](#) by D. A. and André Betz

Comparison to the 4D T-duality chain \Rightarrow Q -flux ?!

T-duality // isometry directions: here 1, 2 $\Rightarrow O(2,2)$;

non-geometry is locally geometric

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Famous toroidal example: NSNS sector:

$T^3 + H_{123} \xleftrightarrow{T_1} \text{twisted torus } (f^1_{23}) \xleftrightarrow{T_2} \text{non-geometric config.}$

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non-geometry is locally geometric

If T-duality // non-isometry: base direction 3

non-geometric, not even locally; characteristic of R -flux

Doubled geometry and Double Field Theory (DFT)

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[hep-th/0406102](#) by C. M. Hull

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Local picture:
$$\begin{array}{ccc} T^2 & \longrightarrow & T^2 \times T^{2'} \\ | & & | \\ S^1 & & S^1 \end{array} \quad T^2, T^{2'} : \text{T-duals}$$

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non-geometry from projection

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Double also the base: doubled geometry

hep-th/0512005 by A. Dabholkar and C. Hull, ... by C. Hull and R. A. Reid-Edwards

$2d$ space with $X^M = (\tilde{x}_m, x^m)$, ∂_M, \tilde{x}_m : dual or winding coord.

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Difficult to construct on concrete examples beyond torus

arXiv:1106.6291 by M. B. Schulz

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Recent work proposes gen. diffeos \mathcal{F}_M^N , can $\rightarrow O(d,d)$

[arXiv:1207.4198](#) by O. Hohm and B. Zwiebach; ... ; [arXiv:1402.2513](#) by M. Cederwall

$$A'_M(X') = \mathcal{F}_M^N A_N(X)$$

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$\mathcal{F}_M{}^N \neq \partial X^N / \partial X'^M \rightarrow$ space: not doubled manifold, η preserved.

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- Asymmetric orbifolds: asymmetric (L/R) orbifold action on \mathcal{X}^m ; no target space description; "*R*-flux situations"?

[hep-th/0210209](#), [hep-th/0512005](#) by A. Dabholkar and C. Hull

[hep-th/0511126](#) by A. Flornoy, B. Williams

[hep-th/0604191](#) by S. Hellerman, J. Walcher

[arXiv:1202.6366](#), [arXiv:1307.0999](#) by C. Condeescu, I. Florakis, C. Kounnas, D. Lüst

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- Closed string non-commutativity: on concrete non-geometric backgrounds: $[\mathcal{X}^m(\tau, \sigma), \mathcal{X}^n(\tau, \sigma)] \neq 0$

[arXiv:1010.1361](#) by D. Lüst

[arXiv:1202.6366](#) by C. Condeescu, I. Florakis, D. Lüst

[arXiv:1211.6437](#) by D. A., M. Larfors, D. Lüst, P. Patalong

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$$[\mathcal{X}^m(\tau, \sigma), \mathcal{X}^n(\tau, \sigma)] \sim N^p Q_p^{mn}$$

\hookrightarrow due to non-geometry, probed by winding

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[arXiv:1010.1263](#), [1106.0316](#), [1112.4611](#) by R. Blumenhagen, et al.

[arXiv:1207.0926](#), [1312.1621](#), [1402.7306](#) by D. Mylonas, P. Schupp, R. Szabo

[arXiv:1309.3172](#) by I. Bakas, D. Lüst

[arXiv:1312.0719](#) by R. Blumenhagen, M. Fuchs, F. Hassler, D. Lüst, R. Sun

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[arXiv:1207.0926](#), [1312.1621](#), [1402.7306](#) by D. Mylonas, P. Schupp, R. Szabo

[arXiv:1309.3172](#) by I. Bakas, D. Lüst

[arXiv:1312.0719](#) by R. Blumenhagen, M. Fuchs, F. Hassler, D. Lüst, R. Sun

Star-products, 3-products (check commutativity, associativity); doubled formalisms

More on closed string non-commutativity

arXiv:1211.6437 by D. A., M. Larfors, D. Lüst, P. Patalong

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$$T^3 + H_{123} \xleftrightarrow{T_1} \text{twisted torus } (f^1_{23}) \xleftrightarrow{T_2} \text{non-geometric config.}$$

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$$[Y^m(\tau, \sigma), Y^n(\tau, \sigma')] = 0$$

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Canonical commutation relations, $\mathcal{P}_m \equiv \frac{\delta \mathcal{L}_\Sigma}{\delta \partial_\tau \mathcal{X}^m}$

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Quantize string Y on twisted torus, using dilute flux approx.,
apply T-duality relations between Y, Z , compute $[Z, Z] \neq 0$

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Doubled phase space interpretation, (also boundary conditions)

$$[Y, Y] = 0 \qquad [Z, Z] \neq 0$$

$$[\tilde{Y}, \tilde{Y}] \neq 0 \qquad [\tilde{Z}, \tilde{Z}] = 0$$

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Doubled phase space interpretation, (also boundary conditions)

$$[Y, Y] = 0 \qquad [Z, Z] \neq 0$$

$$[\tilde{Y}, \tilde{Y}] \neq 0 \qquad [\tilde{Z}, \tilde{Z}] = 0$$

non-geometry/non-commutativity only occurs upon projection

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Earlier results, and related work:

[arXiv:1106.4015](#), [1202.3060](#), [1204.1979](#) by D. A., O. Hohm, M. Larfors, D. Lüst, P. Patalong
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\Leftrightarrow reparametrization of gen. metric \mathcal{H} , i.e. new gen. vielbein

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Apply the field redefinition to $\mathcal{L}_{\text{NSNS}}(g, b, \phi)$

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$$Q_c{}^{ab} = \partial_c \beta^{ab} - 2\beta^{d[a} f^b]{}_{cd}, \quad R^{abc} = 3\beta^{d[a} \nabla_d \beta^{bc]}$$

arXiv:0807.4527 by M. Graña, R. Minasian, M. Petrini, D. Waldram

arXiv:1109.0290 by G. Aldazabal, W. Baron, D. Marqués, C. Núñez

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$$+ 2\eta_{ab} \beta^{ad} \partial_d Q_c{}^{bc} - \eta_{cd} Q_a{}^{ac} Q_b{}^{bd} - \frac{1}{2} \eta_{cd} Q_a{}^{bc} Q_b{}^{ad} - \frac{1}{4} Q^2$$

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$$+ 2\eta_{ab} \beta^{ad} \partial_d Q_c{}^{bc} - \eta_{cd} Q_a{}^{ac} Q_b{}^{bd} - \frac{1}{2} \eta_{cd} Q_a{}^{bc} Q_b{}^{ad} - \frac{1}{4} Q^2$$

10D theory with non-geometric fluxes, uplift 4D \checkmark

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- For standard ∇_m and Levi-Civita connection Γ_{np}^m

$$\tilde{e}^a{}_m \tilde{e}^n{}_b \nabla_n V^m = \nabla_b V^a \equiv \partial_b V^a + \omega_{bc}^a V^c$$

$$\Leftrightarrow \omega_{bc}^a \equiv \tilde{e}^n{}_b \tilde{e}^a{}_m (\partial_n \tilde{e}^m{}_c + \tilde{e}^p{}_c \Gamma_{np}^m)$$

$$\omega_{bc}^a = \frac{1}{2} \left(f^a{}_{bc} + \eta^{ad} \eta_{ce} f^e{}_{db} + \eta^{ad} \eta_{be} f^e{}_{dc} \right)$$

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- There is a new covariant derivative

$$\check{\nabla}^m V^p = -\beta^{mn} \partial_n V^p - \check{\Gamma}_n^{mp} V^n, \quad \check{\nabla}^m V_p = -\beta^{mn} \partial_n V_p + \check{\Gamma}_p^{mn} V_n$$

$$2\check{\Gamma}_p^{mn} = \tilde{g}_{pq} (\beta^{rm} \partial_r \tilde{g}^{nq} + \beta^{rn} \partial_r \tilde{g}^{mq} - \beta^{rq} \partial_r \tilde{g}^{mn}) + 2\tilde{g}_{pq} \tilde{g}^{r(m} \partial_r \beta^{n)q} - \partial_p \beta^{mn}$$

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Proceeding similarly for $\check{\nabla}^m$

$$\tilde{e}^m{}_a \tilde{e}^b{}_n \check{\nabla}^n V_m = \check{\nabla}^b V_a \equiv -\beta^{bd} \partial_d V_a - \omega_{Qa}^{bc} V_c$$

$$\Leftrightarrow -\omega_{Qa}^{bc} \equiv \tilde{e}^b{}_n \tilde{e}^m{}_a \left(-\beta^{nq} \partial_q \tilde{e}^c{}_m + \tilde{e}^c{}_p \check{\Gamma}_{mp}^n \right)$$

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$$\tilde{e}^a{}_m \tilde{e}^n{}_b \nabla_n V^m = \nabla_b V^a \equiv \partial_b V^a + \omega_{bc}^a V^c$$

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$$\omega_{bc}^a = \frac{1}{2} \left(f^a{}_{bc} + \eta^{ad} \eta_{ce} f^e{}_{db} + \eta^{ad} \eta_{be} f^e{}_{dc} \right)$$

- There is a new covariant derivative

$$\check{\nabla}^m V^p = -\beta^{mn} \partial_n V^p - \check{\Gamma}_n^{mp} V^n, \quad \check{\nabla}^m V_p = -\beta^{mn} \partial_n V_p + \check{\Gamma}_p^{mn} V_n$$

$$2\check{\Gamma}_p^{mn} = \tilde{g}_{pq} (\beta^{rm} \partial_r \tilde{g}^{nq} + \beta^{rn} \partial_r \tilde{g}^{mq} - \beta^{rq} \partial_r \tilde{g}^{mn}) + 2\tilde{g}_{pq} \tilde{g}^{r(m} \partial_r \beta^{n)q} - \partial_p \beta^{mn}$$

Proceeding similarly for $\check{\nabla}^m$

$$\tilde{e}^m{}_a \tilde{e}^b{}_n \check{\nabla}^n V_m = \check{\nabla}^b V_a \equiv -\beta^{bd} \partial_d V_a - \omega_{Q_a}{}^{bc} V_c$$

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- For standard ∇_m and Levi-Civita connection Γ_{np}^m

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Define as well a new "Ricci tensor and scalar".

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Define as well a new "Ricci tensor and scalar".

Structures appear clearly using Generalized (Complex)

Geometry formalism; also Lie algebroid structure.

Non-geometry/geometry: global aspects

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Non-geometry/geometry: global aspects

Field redefinition always possible off-shell/locally

$$\mathcal{L}_{\text{NSNS}}(g, b, \phi) = \tilde{\mathcal{L}}_{\beta}(\tilde{g}, \beta, \tilde{\phi}) + \partial(\dots)$$

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Toroidal example :

$$g = f_0 \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & \frac{1}{f_0} \end{pmatrix}, \quad b = f_0 \begin{pmatrix} 0 & -Kz^3 & 0 \\ Kz^3 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}, \quad f_0(z^3) = \frac{1}{1 + (Kz^3)^2}$$

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Geometry restored : T^3 , Q -flux \checkmark .

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Similar for Q -brane \Rightarrow Bianchi identity and Poisson equation.

Non-geometry/geometry: global aspects

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True for a class of non-geometric/geometric backgrounds:
glue with " β -transforms" $\in O(n, n)$ and diffeos.

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Alternative description of 10D non-geometric backgrounds;
compactification possible.

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Different perspectives on non-geometry:

- 4D non-geometric fluxes in gauged supergravity
- 10D non-geometric backgrounds (stringy symmetry); toroidal example, Q -brane
- T-fold, doubled geometry, and DFT
- Generalized Complex Geometry and Gen. tangent bundle
- World-sheet: asymmetric orbifolds, non-commutativity and non-associativity

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β -supergravity: 10D theory with non-geometric fluxes Q, R
non-geometric bckgd of standard supergravity \rightarrow geometric
 \leftrightarrow compactification, uplift of 4D gauged supergravity

More about it, Generalized Geometry, NS -branes: talk 27/03

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More about it, Generalized Geometry, NS -branes: talk 27/03

Study various geometric structures in different perspectives.

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Different perspectives on non-geometry:

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More about it, Generalized Geometry, NS -branes: talk 27/03

Study various geometric structures in different perspectives.

Extension beyond the NSNS sector: RR non-geometric fluxes?
exotic D -branes? Exceptional geometry/field theory could help
SUSY with non-geometric fluxes, $SU(3) \times SU(3)$ structures
 \leftrightarrow Get new pheno. interesting backgrounds, de Sitter vacua.